

The Center for Astrophysical Thermonuclear Flashes

Various Hydro Solvers in FLASH3

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FLASH3 Tutorial
June 22-23, 2009



An Advanced Simulation and Computing (ASC)
Academic Strategic Alliances Program (ASAP)
Center at The University of Chicago





Hydro Solvers in FLASH3



- ❑ In FLASH3, “Hydro” unit houses more than one usual gas dynamics solver:
 - ❑ Pure hydrodynamics (i.e., gas dynamics) solvers (PPM & MUSCL-Hancock)
 - ❑ Magnetohydrodynamics(MHD) solvers (Unsplit Staggered Mesh & 8-wave)
 - ❑ Relativistic hydrodynamics (RHD) solver

- ❑ The Hydro unit is organized into two different subunits depending on how you treat multidimensional flux updates:
 - ❑ Operator (dimensional) Splitting (Strang, 1968) vs. Unsplit (Colella, 1990; Lee & Deane, 2009)
 - ❑ `source/Hydro/HydroMain/split` (PPM, 8-wave MHD, RHD)
 - ❑ `source/Hydro/HydroMain/unsplit` (Staggered Mesh MHD, MUSCL-Hancock pure-Hydro)

- ❑ All these five major different solvers are based on high-order Godunov (1959) method which involves:
 - ❑ Finite volume method
 - ❑ Predictor-corrector
 - ❑ Riemann problem
 - ❑ Explicit time advancement



Physics of Hydro Solvers



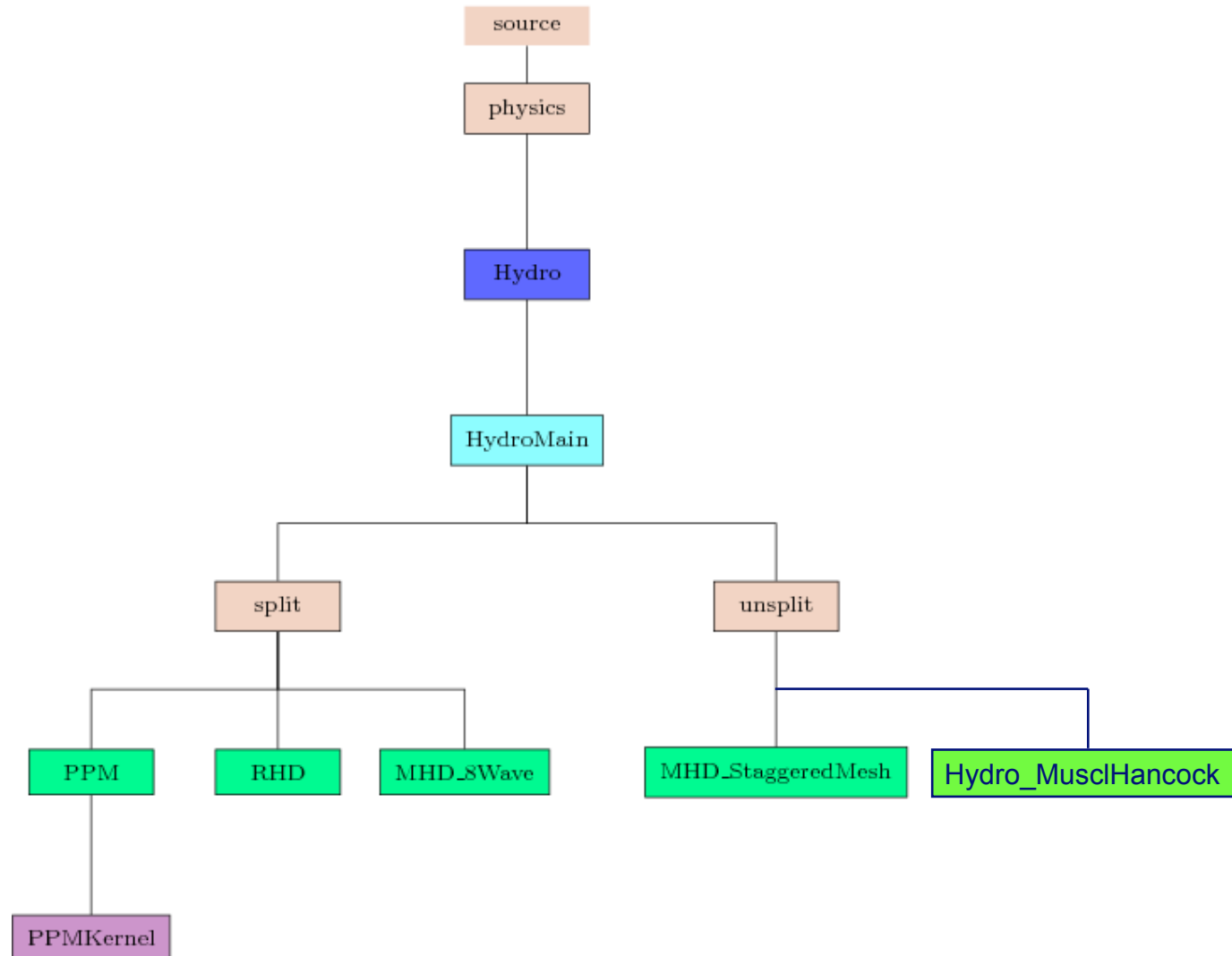
- ❑ Pure hydrodynamics solvers (PPM & MUSCL-Hancock)
 - ❑ Compressible reactive gas dynamics
 - ❑ Can solve a broad range of (astro)physical problems

- ❑ MHD solvers (Unsplit Staggered Mesh & 8-wave)
 - ❑ flows of conducting fluids (ionized gases, liquid metals) in presence of magnetic fields
 - ❑ Plasma is a completely ionized gas, consisting of freely moving positively charged ions (or nuclei) and negatively charged electrons
 - ❑ Lorentz forces act on charged particles and change their momentum and energy. In return, particles alter strength and topology of magnetic fields.
 - ❑ A valid macroscopic model of magnetized plasma → MHD

- ❑ Relativistic hydrodynamics solver (RHD)
 - ❑ A wide variety of astrophysical flows exhibit relativistic behavior
 - ❑ accretion around compact objects, jets in extragalactic radio sources, pulsar winds, gamma ray bursts



Hydro Unit in FLASH3

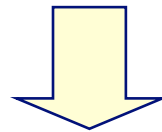




Operator Splitting vs. Unsplit Formulations



$$\frac{\partial U}{\partial t} + \nabla \cdot \text{Flux} = 0$$



$$\frac{\partial U}{\partial t} + \frac{\partial F}{\partial x} + \frac{\partial G}{\partial y} = 0$$



Operator Splitting vs. Unsplit Formulations



$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} + \frac{\partial \mathbf{G}}{\partial y} = \mathbf{0}$$

Splitting

$$\mathbf{X}^{\Delta t} : \left. \begin{array}{l} \text{PDE} : U_t + \mathbf{F}(\mathbf{U})_x = 0 \\ \text{IC} : U^n \end{array} \right\} \xrightarrow{\Delta t} U^{n+1/2}$$

$$\mathbf{Y}^{\Delta t} : \left. \begin{array}{l} \text{PDE} : U_t + \mathbf{G}(\mathbf{U})_y = 0 \\ \text{IC} : U^{n+1/2} \end{array} \right\} \xrightarrow{\Delta t} U^{n+1}$$

1st order Strang Splitting

$$U^{n+1} = \mathbf{X}^{\Delta t} \mathbf{Y}^{\Delta t} U^n$$

2nd order Strang Splitting

$$U^{n+1} = \left(\mathbf{X}^{\Delta t/2} \mathbf{Y}^{\Delta t/2} \right) \left(\mathbf{Y}^{\Delta t/2} \mathbf{X}^{\Delta t/2} \right) U^n$$

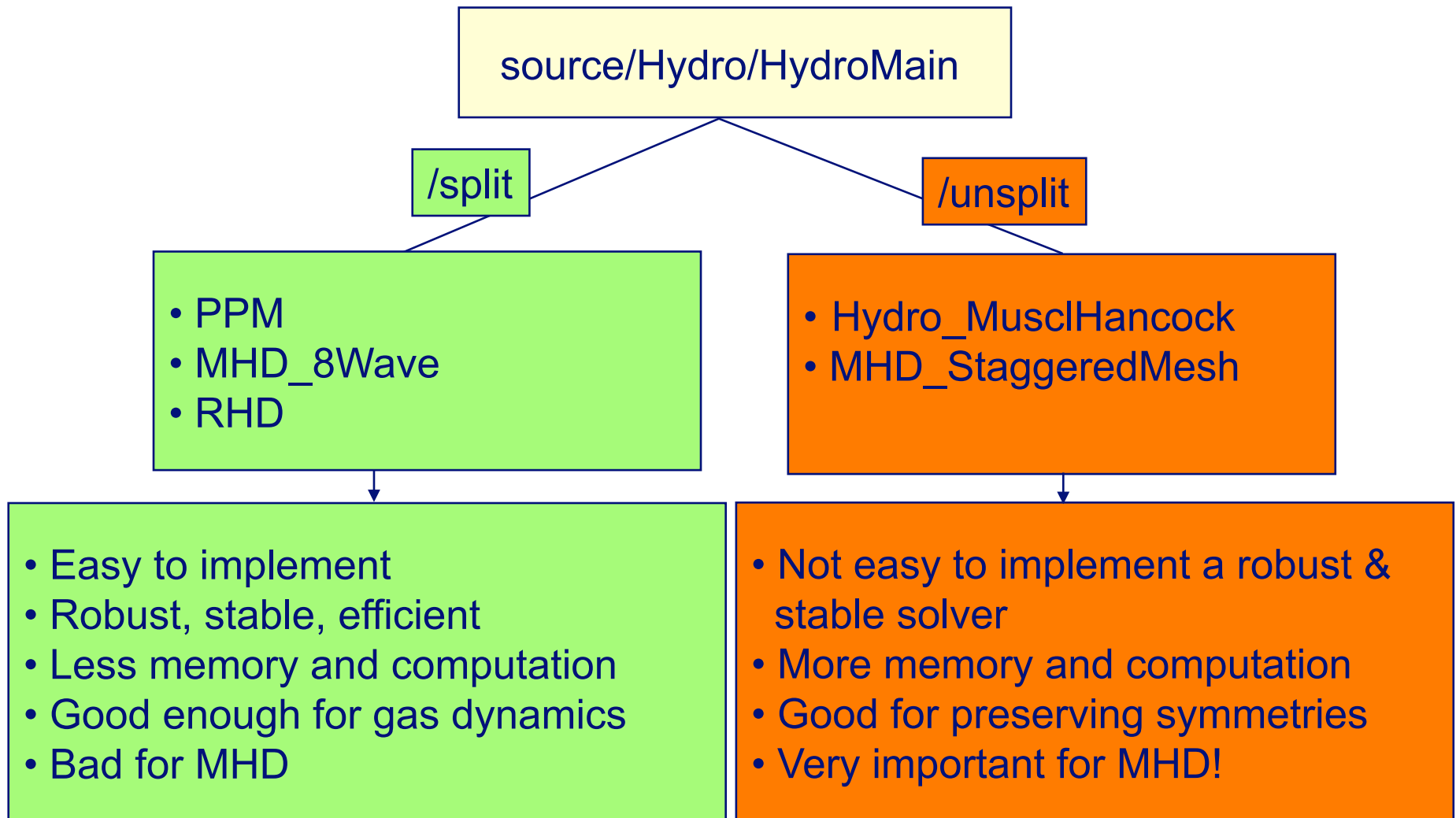
Unsplit

$$\text{PDE} : U_t + \mathbf{F}(\mathbf{U})_x + \mathbf{G}(\mathbf{U})_y = 0$$

$$\text{IC} : U(x, y, t^n) = U^n$$



Splitting vs. Unsplit

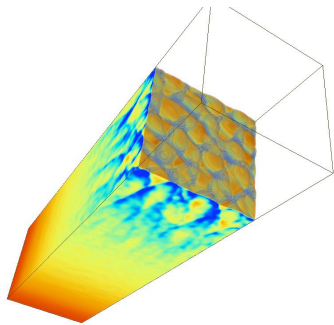




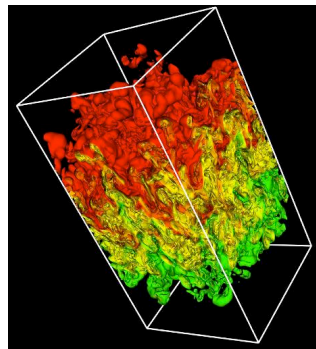
Hydro Solvers Primer – Split solvers



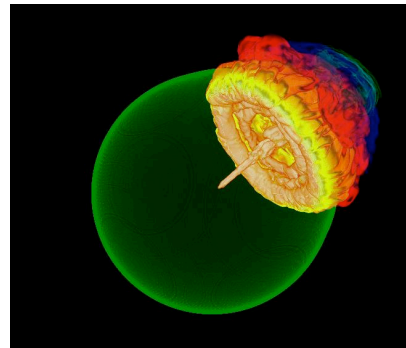
- ❑ **Piecewise-parabolic method solver (PPM) (Fryxell et al., 2000)**
 - ❑ Parabolic interpolation of data over each cell (Colella and Woodward, 1984)
 - ❑ (Ideally) 3rd order, (formally) 2nd order, (practically) 1st order in shocks and discontinuities in spatial discretization
 - ❑ High resolution with accuracy (smooth flows)
 - ❑ Monotonocity enforcement, interpolant flattening, steepening of contact discontinuities
 - ❑ 2nd order in explicit time evolution using operator splitting formulation



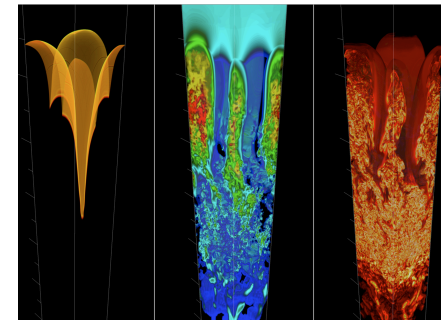
Cellular detonation



Rayleigh-Taylor instability



Gravitationally confined detonation



Turbulent Nuclear Burning



Hydro Solvers Primer – Split solvers

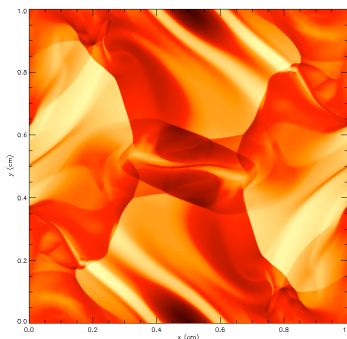


❑ MHD 8-wave solver (Timur Linde, 1999)

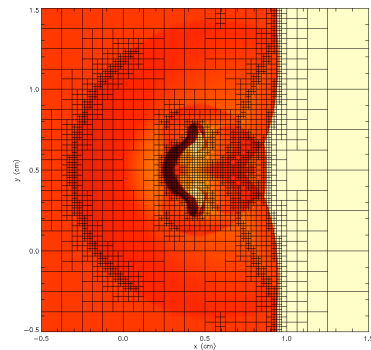
❑ Monotone Upstream-centered Scheme for Conservation Laws (MUSCL) approach (Van Leer, 1977)

- ❑ 2nd order in space, 2nd order in time
- ❑ Magnetic monopoles (8th wave) are convected away, rather than accumulated (Powell et al., 1998) (i.e., $\nabla \cdot \mathbf{B} \neq 0$)
- ❑ Non-conservative formulation of the MHD governing equations
- ❑ Incorrect jump conditions and incorrect propagation speeds across discontinuities
- ❑ Robust and accurate (as compared to the basic conservative scheme)

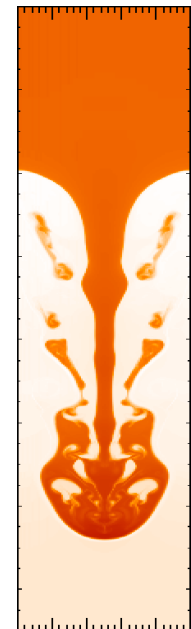
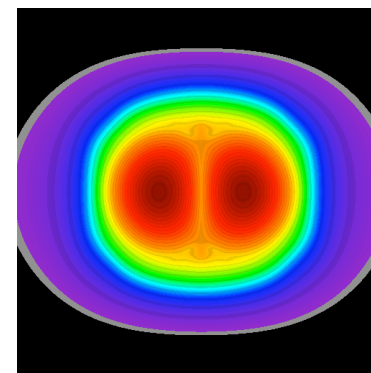
Orszag-Tang



Shock-Cloud Interaction



Magnetic reconnection

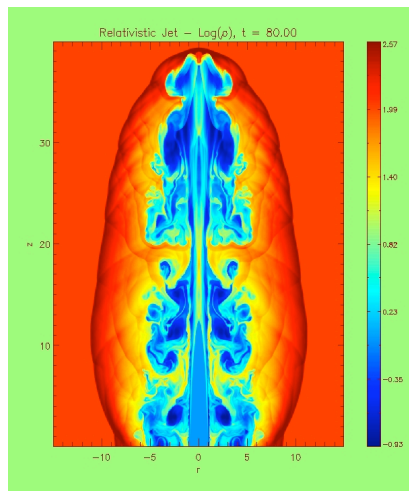




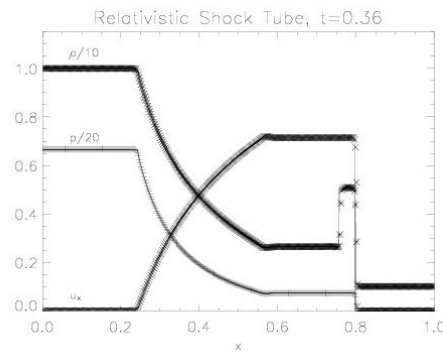
Hydro Solvers Primer – Split solvers



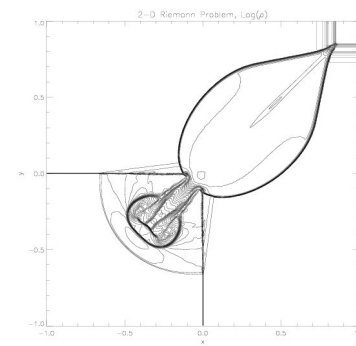
- ❑ **Special Relativistic solver (RHD) (A. Mignone, 2004)**
 - ❑ PPM (3rd order in space) and PLM (2nd order) interpolations
 - ❑ 2nd order in explicit time evolution using operator splitting formulation
 - ❑ (special) relativistic effects are twofold:
 - kinematical, $v \sim c$ ($\gamma = 1/(1 - v^2)^{1/2} \gg 1$)
 - thermodynamical, $c_s \sim c$
 - ❑ Relativistic flows with $\gamma > (3/2)^{1/2}$ are always supersonic, and therefore shock-capturing methods are essential (Martí and Müller, 2003)



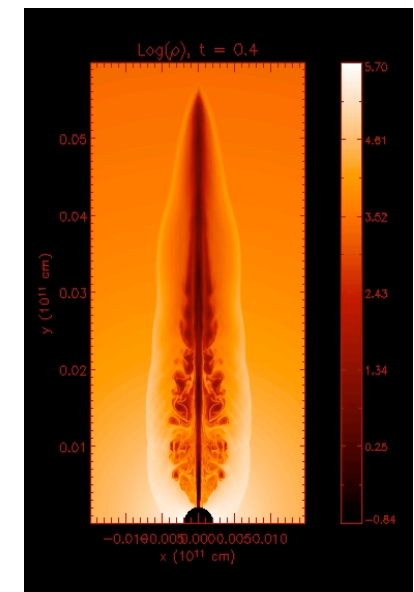
$\gamma = 10$ Jet



Relativistic Shock tube



2-D Riemann Problem



Jet through collapsars (GRB), $\gamma \sim 50$

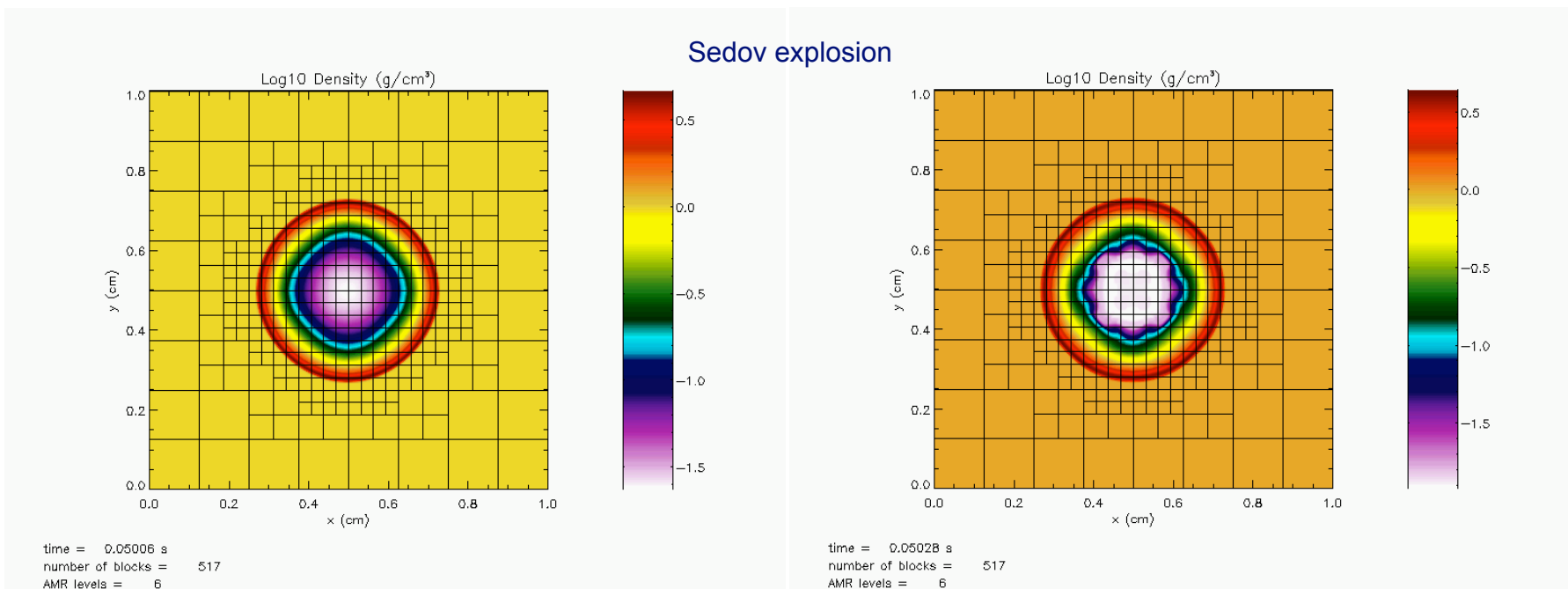


Hydro Solvers Primer – Unsplit solvers



Unsplit pure-Hydro solver (Lee, 2009)

- Reduced version of USM-MHD solver without magnetic and electric fields
- 2nd order MUSCL-Hancock in space and time
- Preserves better flow symmetries (Roe solver + Carbuncle instability fix)



Unsplit MUSCL-Hancock

Split PPM

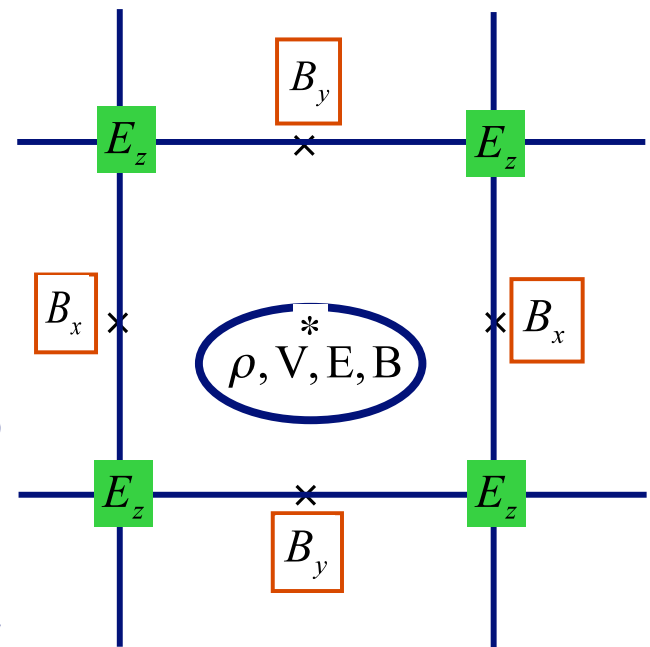


Hydro Solvers Primer – Unsplit solvers



Unsplit Staggered Mesh (USM) MHD solver (Lee, 2006; Lee and Deane 2009)

- A very efficient new data reconstruction algorithm
 - MUSCL-Hancock (2nd order in space) type characteristic tracing method
 - No extra Riemann problems for transverse fluxes!
- A new way of treating multidimensional MHD source terms in unsplit formulation
- Constrained Transport (CT) algorithm (Evans and Hawley, 1988; Balsara and Spicer, 1999) for induction equations to maintain $\nabla \cdot \mathbf{B} = 0$ on a staggered grid
 - cell-centered, face-centered, corner (edge-centered) variables
- Enhanced solution accuracy in calculating electric fields for the induction equations (modified electric field construction)
 - Added proper amount of numerical dissipation – important!



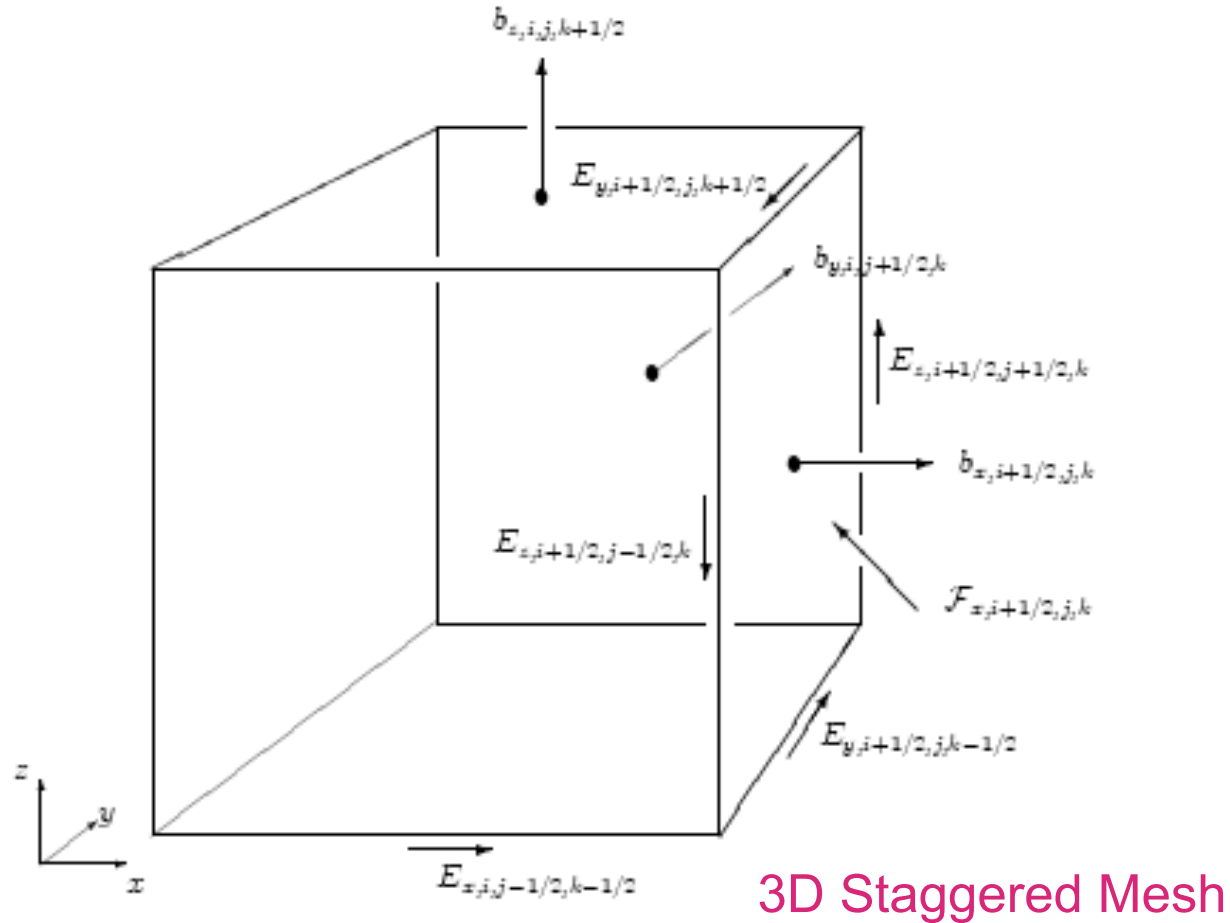
2D staggered mesh



Hydro Solvers Primer – Unsplit solvers



- Unsplit Staggered Mesh (USM) MHD solver (Lee, 2006; Lee and Deane 2009)



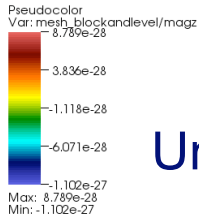


Why unsplit formulation is necessary for MHD?

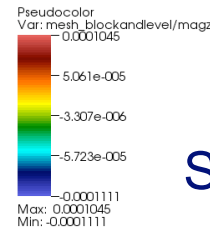
- Operator splitting MHD schemes cannot avoid erroneous growth in B_z :

$$\frac{\partial B_z}{\partial t} + B_z \frac{\partial u}{\partial x} - B_x \frac{\partial w}{\partial x} - w \frac{\partial B_x}{\partial x} + u \frac{\partial B_z}{\partial x} + B_z \frac{\partial v}{\partial y} - B_y \frac{\partial w}{\partial y} - w \frac{\partial B_y}{\partial y} + v \frac{\partial B_z}{\partial y} = 0$$

$$w \nabla \cdot \mathbf{B} = w (\Delta B_{x,i} / \Delta x + \Delta B_{y,j} / \Delta y)$$

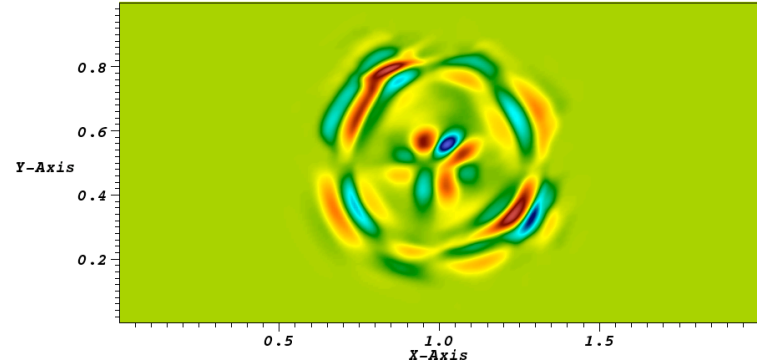
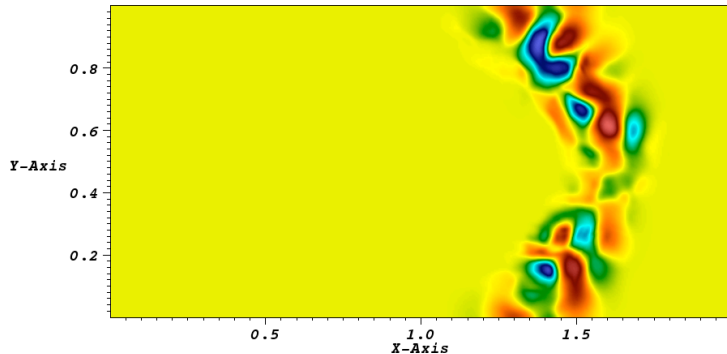


Unsplit Staggered Mesh solver



Split 8-wave solver

2D Field Loop advection test





More on the USM-MHD solver



- ❑ Physics
 - ❑ Ideal and non-ideal flows
 - ❑ Magnetic resistivity, thermal conductivity, and viscosity
 - ❑ EOS
 - ❑ Ideal gamma, multiple gamma, Helmholtz (degenerate EOS)
 - ❑ Gravity
 - ❑ Multiple species, particles
 - ❑ Well tested for wide ranges of plasma flows: $10^{-6} < \beta (\equiv p/B_p) < 10^6$

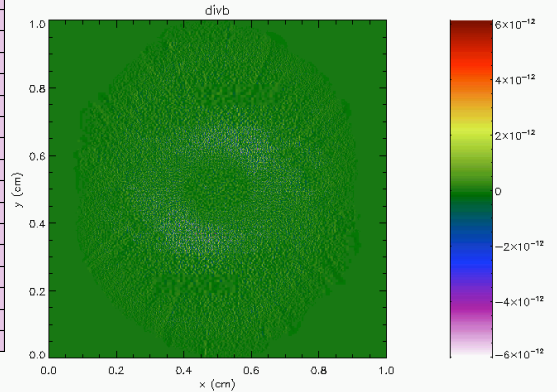
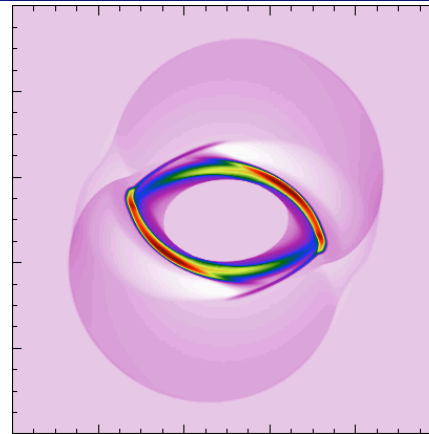
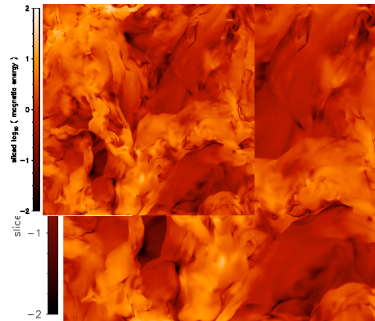
- ❑ Implementations and algorithms
 - ❑ Riemann solvers
 - ❑ Roe (default), HLLC, HLLD (robust and accurate, suggested for most plasma flows)
 - ❑ Carbuncle, even-odd instability fix for Roe solver
 - ❑ Strong shock-rarefaction detect algorithm (Balsara)
 - ❑ Various slope limiters (Minmod, MC, Van Leer, hybrid)
 - ❑ Two prolongation methods of divergence-free B fields on AMR
 - ❑ Use of face-centered variables, and edge-centered variable
 - ❑ Wide ranges of CFL limit: CFL < 1 for 1D, 2D and 3D



Applications

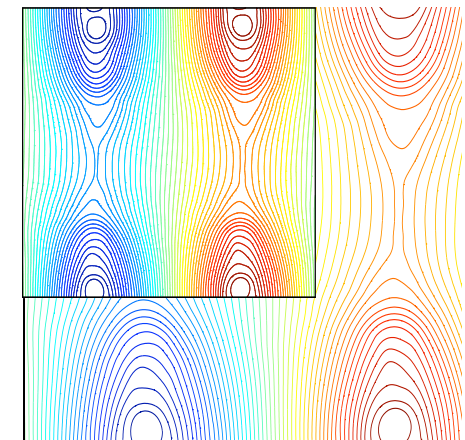


High Mach Number MHD Turbulence

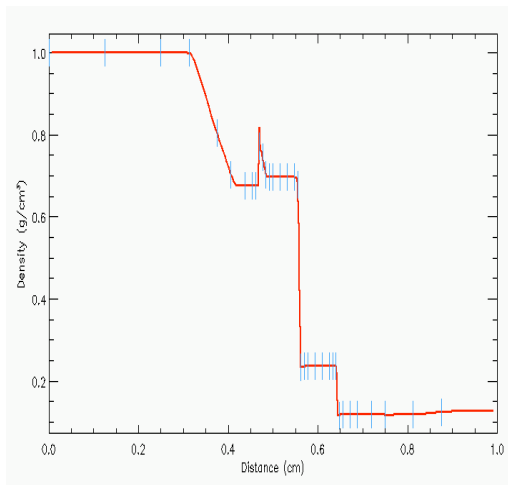


time = 0.15035 a
number of blocks = 2461
AMR levels = 7

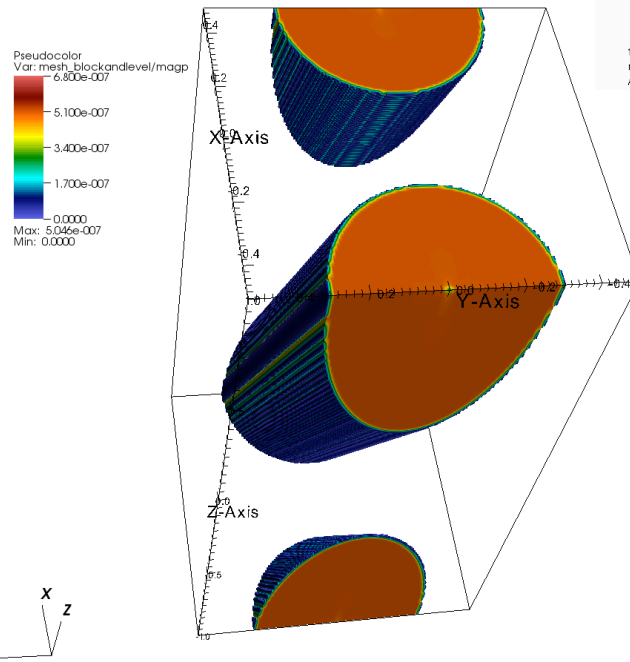
Rotor



Magnetic Reconnection



Brio-Wu MHD shock-tube



3D Field Loop advection



How to Setup a New Problem



- ❑ MHD Simulation (especially with the USM solver) should be located in
 - ❑ source/Simulation/SimulationMain/magnetoHD/
 - ❑ Special prolongation for face-centered variables

- ❑ Create Simulation_initBlock.F90 exactly as you would do for hydro. Just do not forget to set magnetic field variables (both cell-centered and cell face-centered) in the initialization routine.
 - ❑ Magnetic fields need to satisfy $\nabla \cdot \mathbf{B} = 0$

- ❑ Do not add magnetic pressure to total specific energy, because FLASH EOS routines assume a specific expression for it.

- ❑ Create Config and flash.par files for your own Simulation directory.

- ❑ Special care in writing custom boundary conditions in Grid_bcApplyToRegionSpecialized.F90.

- ❑ Write custom functions and do not forget to add them to Makefile. Such custom functions in your Simulation directory will override other standard implementations.



Example – Config



```

! Configuration file for Orszag Tang MHD vortex problem
# (Orszag and Tang, J. Fluid Mech., 90:129--143, 1979)

REQUIRES physics/Hydro/HydroMain
REQUIRES physics/Eos/EosMain/Gamma

USESETUPVARS withParticles

IF withParticles
  PARTICLETYPE passive INITMETHOD lattice MAPMETHOD quadratic

  REQUIRES Particles/ParticlesMain
  REQUESTS IO/IOMain
  REQUESTS IO/IOParticles
  REQUESTS Particles/ParticlesMapping/Quadratic
  REQUESTS Particles/ParticlesInitialization/Lattice
ENDIF

D tiny Threshold value used for numerical zero
PARAMETER tiny REAL 1.e-16

# ----- For Resistive MHD setup -----#
#REQUIRES physics/materialProperties/Conductivity/ConductivityMain/Constant-diff
#REQUIRES physics/materialProperties/Viscosity/ViscosityMain
#REQUIRES physics/materialProperties/MagneticResistivity/MagneticResistivityMain
#REQUIRES physics/sourceTerms/Diffuse/DiffuseMain

#VARIABLE vecz # vector potential Az
#----- End of Resistive MHD setup -----#
```



Example – flash.par



```
#      DivB control switch
killdivb      = .true.

#      Flux Conservation for AMR
flux_correct  = .true.

## -----##
## SWITCHES SPECIFIC TO THE UNSPLIT STAGGERED MESH MHD SOLVER ##
#      I. INTERPOLATION SCHEME:
order         = 2      # Interpolation order (First/Second order)
slopeLimiter  = "mc"   # Slope limiters (mirmod, mc, vanLeer, hybrid, limited)
LimitedSlopeBeta= 1.   # Slope parameter for the "limited" slope by Toro
charLimiting  = .true. # Characteristic limiting vs. Primitive limiting

#      II. MAGNETIC(B) AND ELECTRIC(E) FIELDS:
E_modification = .true. # High order algorithm for E-field construction
energyFix      = .true. # Update magnetic energy using staggered B-fields
ForceHydroLimit = .false. # Pure Hydro Limit (B=0)
prolMethod     = "injection_prol" # Prolongation method (injecton_prol, balsara_prol)

#      III. RIEMANN SOLVERS:
RiemannSolver  = "Roe" # Roe, HLL, HLLC, HLLD, LF
shockInstabilityFix = .false. # Carbuncle instability fix for the Roe solver
entropy        = .false. # Entropy fix for the Roe solver

#      IV. STRONG SHOCK HANDELING SCHEME:
shockDetect    = .false. # Shock Detect for numerical stability
## -----##
```



Example – setup lines



- ❑ `./setup magnetoHD/OrszagTang -auto -2d +usm (+8wave) -opt (-debug) -objdir=OT2D -with-unit=Particles +pm3 (+pm4dev) -nxb=8 -nyb=8`
 - ❑ `lrefine_min = 1, lrefine_max = 6`

- ❑ `./setup magnetoHD/Rotor -auto -2d +usm -opt +ug -nxb=200 -nyb=300`
 - ❑ `iProcs = 2, jProcs = 2`

- ❑ `./setup magnetoHD/Rotor -auto -2d +usm -opt +nofbs`
 - ❑ `iGridSize = 400, jGridSize = 600, iProcs = 2, jProcs = 2`

- ❑ More shortcuts can be found in
 - ❑ `/bin/setup_shortcuts.txt`
 - ❑ Users can add their own customized shortcut(s) by editing the file



Example – setup lines



/bin/setup_shortcuts.txt

```
# Choice of Grid
grid:-unit=Grid:
ug:+grid:Grid=UG:
pm2:+grid:Grid=PM2:
pm40:+grid:Grid=PM40:
pm3:+pm40
pm4dev:+grid:Grid=PM4DEV:ParameshLibraryMode=True

# Choice of MHD solver
# NOTE: The 8wave mhd solver only works with the native interpolation.
8wave:--with-unit=physics/Hydro/HydroMain/split/MHD_8Wave:+grid:-gridinterpolation=native

# NOTE: If pure hydro mode used with the USM solver, add +pureHydro in setup
usm:--with-unit=physics/Hydro/HydroMain/unsplit/MHD_StaggeredMesh:--without-unit=physics/Hydro/HydroMain/split/MHD_8Wave
pureHydro:physicsMode=hydro
unsplitHydro:--with-unit=physics/Hydro/HydroMain/unsplit/Hydro_MusclHancock
```



Tips on setting up your problems



- ❑ Simulation's Config file contains all default runtime parameter values specific to your problem
- ❑ default.par is also generated in your object directory
- ❑ Flash.h, setup_units, setup_vars, setup_params, etc
- ❑ Always helpful to read FLASH user's guide
- ❑ Further read/refer references to understand the various roles/effects on using different choices of runtime parameters
 - ❑ e.g., Roe vs. HLL-type Riemann solvers
- ❑ Use “-debug” for testing
- ❑ Many (simple) issues come from wrong initializations and wrong boundary conditions
 - ❑ Simplify your issues as much as possible
 - ❑ Detailed description on your issues/bugs is always welcome
 - ❑ Your simulation files, flash.par, log files, boundary conditions, compilers, setup lines, etc.



Tips in general



- ❑ Make sure you know your problem
 - ❑ Literature research
 - ❑ IC, BC, units
 - ❑ Any working example?
 - ❑ Working on FLASH3

- ❑ Comments, comments, and comments – keep your journal

- ❑ Useful books and references:
 - ❑ LeVeque, Finite volume methods for hyperbolic problems
 - ❑ Toro, Riemann solvers and numerical methods for fluid dynamics
 - ❑ Laney, Computational gas dynamics
 - ❑ Goedbloed and Poedts, Principles of Magnetohydrodynamics
 - ❑ NRL plasma formulary
 - ❑ FLASH user's guide

- ❑ Ask around good people!



Supplementary slides





Ideal MHD Equations



$$\frac{\partial}{\partial t} \begin{pmatrix} \rho \\ \mathbf{V} \\ \rho E \\ \mathbf{B} \end{pmatrix} + \nabla \cdot \begin{pmatrix} \rho \mathbf{V} \\ \rho \mathbf{V} \mathbf{V} + (p + \frac{B^2}{2}) \bar{\mathbf{I}} - \mathbf{B} \mathbf{B} \\ \mathbf{V} (\rho E + p + \frac{B^2}{2}) - \mathbf{B} (\mathbf{V} \cdot \mathbf{B}) \\ \mathbf{V} \mathbf{B} - \mathbf{B} \mathbf{V} \end{pmatrix} = 0$$

Major Properties:

- ❑ MHD equations form a hyperbolic system → Seven families of waves (entropy, Alfvén and fast and slow magnetoacoustic waves).
- ❑ Convex space of physically admissible variables if convex EOS.
- ❑ Non-convex flux function → Multiple degeneracies in the eigensystem, possibility of compound waves, shock evolutionarity concerns.

♪ **Important to remember:** Fluid (Euler) equations are not the limiting case of MHD equations in the $B \rightarrow 0$ case in strict mathematical sense.

♪



MHD Equations in FLASH3



$$\begin{aligned}\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0 \\ \frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla p_* &= \rho \mathbf{g} + \nabla \cdot \tau \\ \frac{\partial \rho E}{\partial t} + \nabla \cdot (\mathbf{v}(\rho E + p_*) - \mathbf{B}(\mathbf{v} \cdot \mathbf{B})) &= \rho \mathbf{g} \cdot \mathbf{v} + \nabla \cdot (\mathbf{v} \cdot \tau + \sigma \nabla T) + \nabla \cdot (\mathbf{B} \times (\eta \nabla \times \mathbf{B})) \\ \frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v}) &= -\nabla \times (\eta \nabla \times \mathbf{B})\end{aligned}$$

where

$$\begin{aligned}p_* &= p + \frac{B^2}{2}, \\ E &= \frac{1}{2} v^2 + \epsilon + \frac{1}{2} \frac{B^2}{\rho}, \\ \tau &= \mu \left((\nabla \mathbf{v}) + (\nabla \mathbf{v})^T - \frac{2}{3} (\nabla \cdot \mathbf{v}) \right)\end{aligned}$$

- ❑ Advective terms are discretized using slope-limited TVD scheme.
- ❑ Diffusive terms are discretized using central finite differences.
- ❑ Time integration is done using one-stage Hancock scheme.
- ❑ Directions are either split or unsplit.



Overview on Numerical MHD codes



- ❑ Numerical MHD
 - ❑ Finite difference method (e.g., ZEUS code by Stone & Norman, 1992)
 - ❑ High order Godunov method – high resolution shock capturing approach

- ❑ Two major algorithmic progresses in high order Godunov based MHD codes
 - ❑ Multidimensional integration algorithm
 - ❑ split vs. unsplit
 - ❑ Donor cell method
 - ❑ Corner transport upwind (CTU), Colella 1990

 - ❑ Divergence-free constraint on B
 - ❑ Hodge projection
 - ❑ Zachary et al, 1994; Crockett et al, 2005
 - ❑ Constrained transport (CT)
 - ❑ Evans & Hawley, 1988; Dai & Woodward, 1998; Ryu et al, 1998; Balsara & Spicer, 1999; Toth, 2000; Londrillo & Del Zanna, 2004; Ziegler, 2004; Gardiner & Stone, 2005; Fromang, 2006; Cunningham et al, 2007
 - ❑ Non-conservative formulation
 - ❑ Powell (et al), 1994; 1999; Falle et al, 1998; Dedner et al, 2002



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 - ❑ Zachary et al, 1994; Crockett et al, 2005
 - ❑ **Constrained transport (CT)**
 - ❑ Evans & Hawley, 1988; Dai & Woodward, 1998; Ryu et al, 1998; Balsara & Spicer, 1999; Toth, 2000; Londrillo & Del Zanna, 2004; Ziegler, 2004; Gardiner & Stone, 2005; Fromang, 2006; Cunningham et al, 2007
 - ❑ Non-conservative formulation
 - ❑ Powell (et al), 1994; 1999; Falle et al, 1998; Dedner et al, 2002

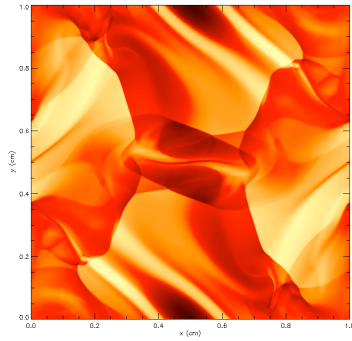
USM scheme



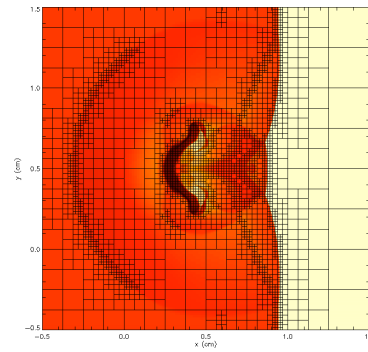
Applications



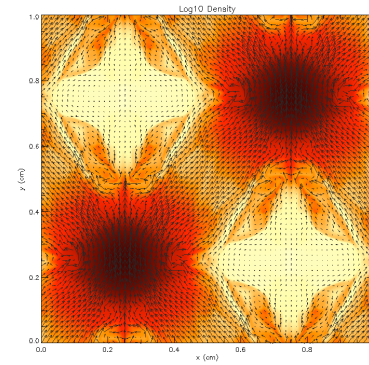
Orszag-Tang



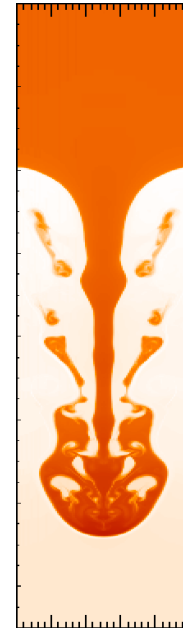
Shock-Cloud Interaction



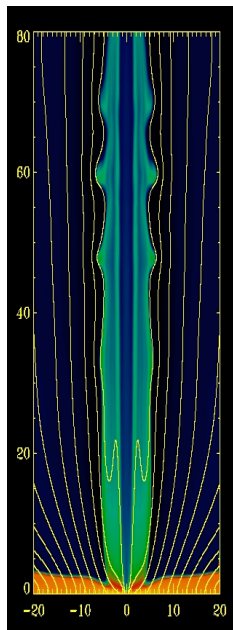
Self-Gravitating Plasma



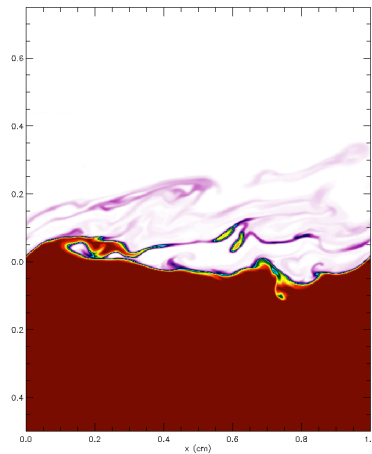
Magnetic RT



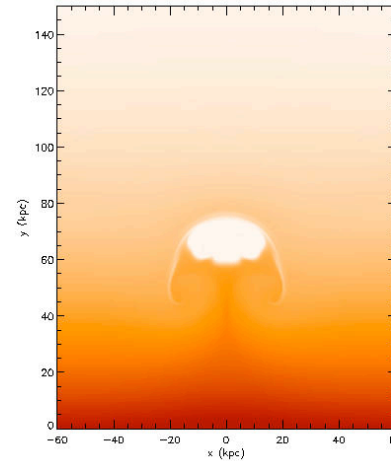
Jet Launching



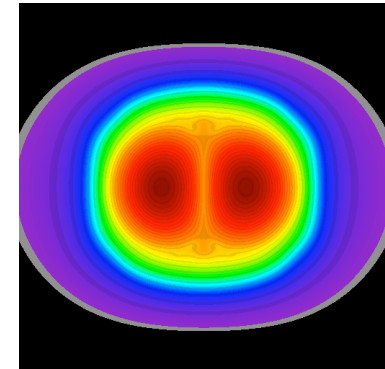
Surface Gravity Wave



Rising bubble



Magnetic reconnection





Beyond Plain MHD



Plasma effects

- ❑ Reduced 2D Hall (Grasso et al, 1999)
- ❑ Electron inertia and compressibility
- ❑ 3D Hall MHD and two-fluid MHD

$$\mathbf{E} + \mathbf{V} \times \mathbf{B} = \frac{1}{S} \mathbf{J} + \frac{d_e^2}{n} \frac{d\mathbf{J}}{dt} + \frac{d_i}{n} (\mathbf{J} \times \mathbf{B} - \nabla \cdot \vec{p}_e)$$

$$\frac{\partial F}{\partial t} + [\phi, F] = \rho_s^2 [U, \psi]$$

$$\frac{\partial U}{\partial t} + [\phi, U] = [J, \psi]$$

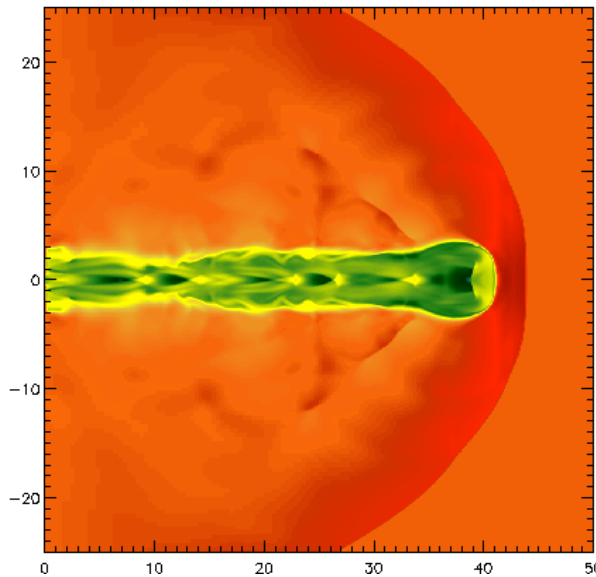
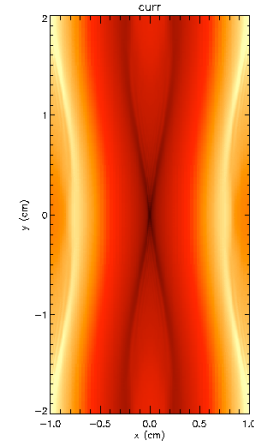
$$F = \psi + d_e^2 J$$

$$J = -\nabla^2 \psi$$

$$U = \nabla^2 \phi$$

$$\vec{B} = B_0 \hat{z} + \nabla \psi \times \hat{z}$$

$$\vec{v} = \hat{z} \times \nabla \phi$$



Relativistic MHD

$$\frac{\partial \mathbf{W}}{\partial t} + (\nabla \cdot \mathbf{F})^T = \mathbf{0}$$

$$\mathbf{W} = \begin{pmatrix} \Gamma \rho \\ \Gamma^2 \frac{e+p}{c^2} \mathbf{u} + \frac{1}{c^2} \mathbf{S}_\Lambda \\ \mathbf{B} \\ \Gamma^2 (e+p) - p - \Gamma \rho c^2 + e_\Lambda \end{pmatrix} \quad \mathbf{F} = \begin{pmatrix} \Gamma \rho \mathbf{u} \\ \frac{\Gamma^2}{c^2} (e+p) \mathbf{u} \mathbf{u} + p \mathbf{I} + \mathbf{P}_\Lambda \\ \mathbf{u} \mathbf{B} - \mathbf{B} \mathbf{u} \\ [\Gamma^2 (e+p) - \Gamma \rho c^2] \mathbf{u} + \mathbf{S}_\Lambda \end{pmatrix}^T$$

$$\Gamma = \frac{1}{\sqrt{1 - \frac{u^2}{c^2}}}$$

$$e_\Lambda = \frac{1}{2\mu_0} \left(B^2 + \frac{1}{c^2} E^2 \right)$$

$$\mathbf{S}_\Lambda = \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B}),$$

$$\mathbf{P}_\Lambda = e_\Lambda \mathbf{I} - \frac{1}{\mu_0} \mathbf{B} \mathbf{B} - \frac{1}{\mu_0 c^2} \mathbf{E} \mathbf{E}.$$

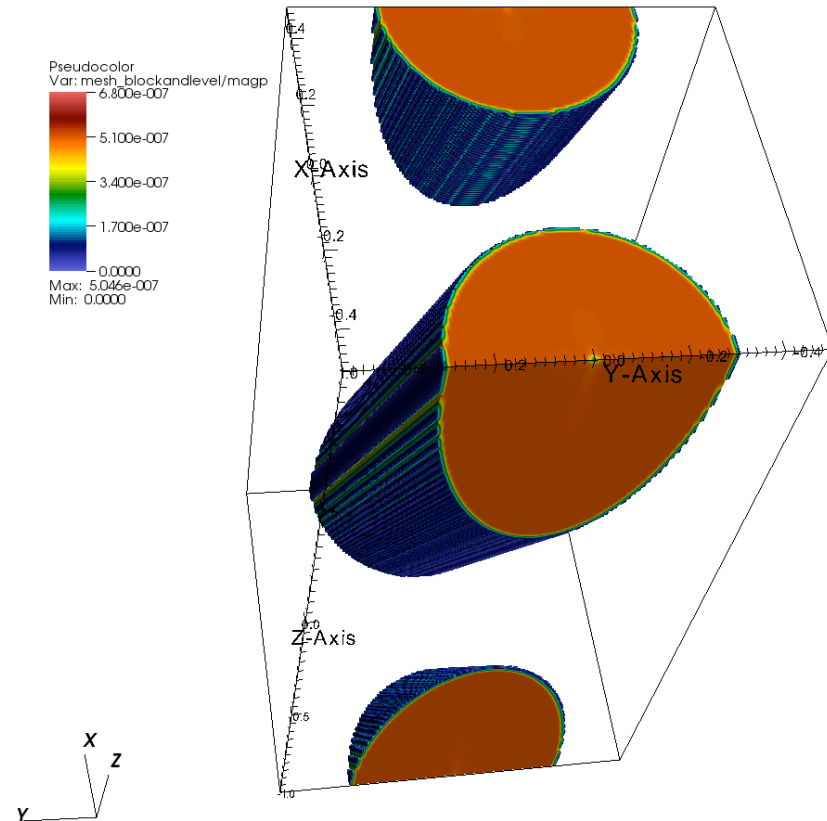
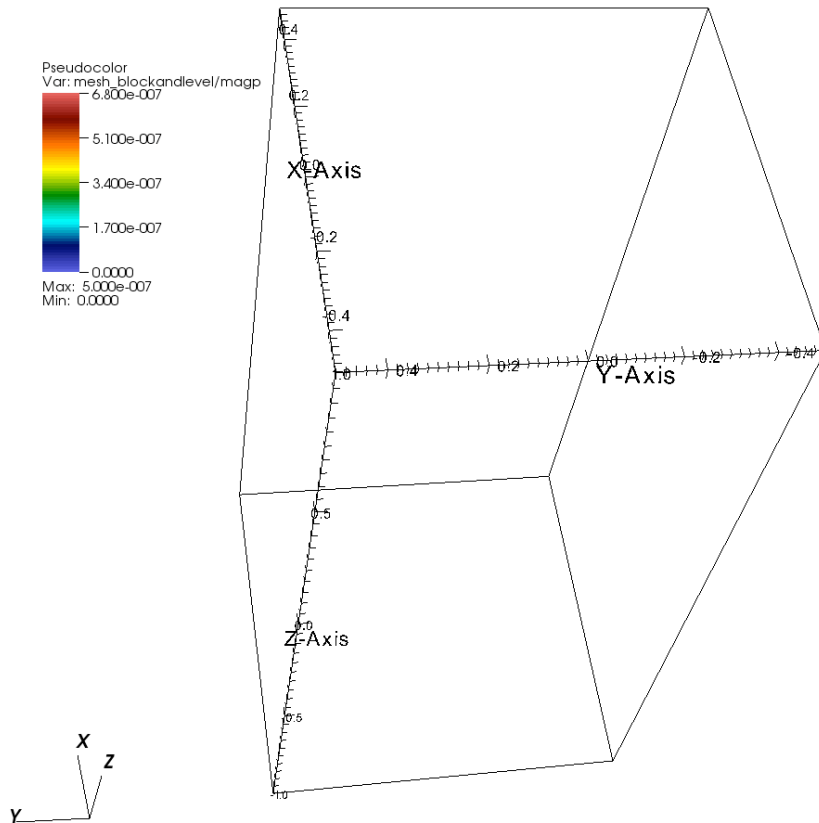


Numerical Results on Benchmarked problems (1)



Field Loop advection in 3D

Gardiner & Stone (2008)



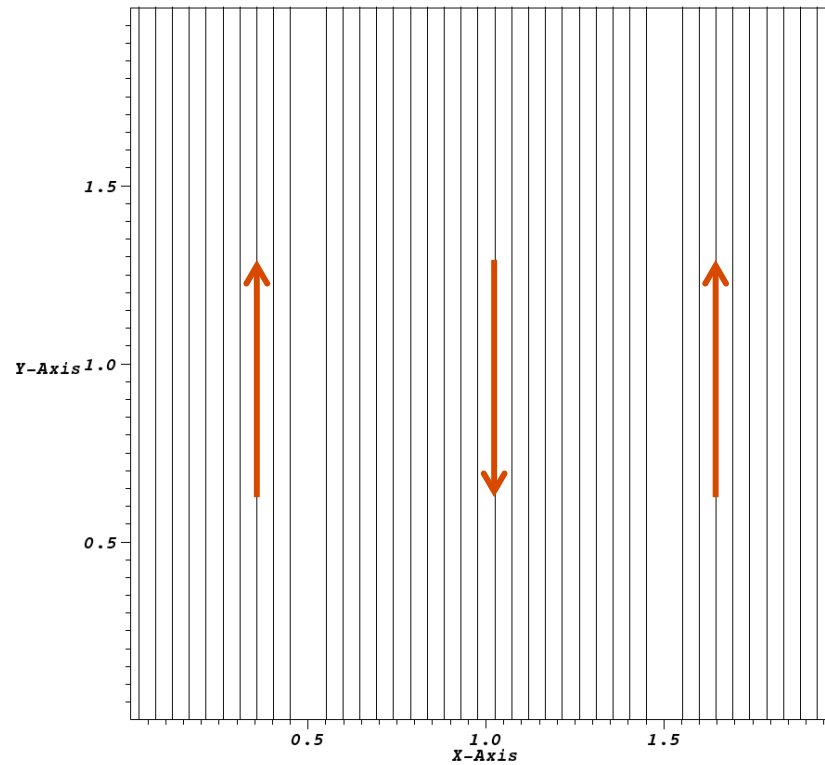


Numerical Results on Benchmarked problems (2)



Current Sheet Problem with different plasma beta values (magnetic field lines are shown)

$$t = 0$$

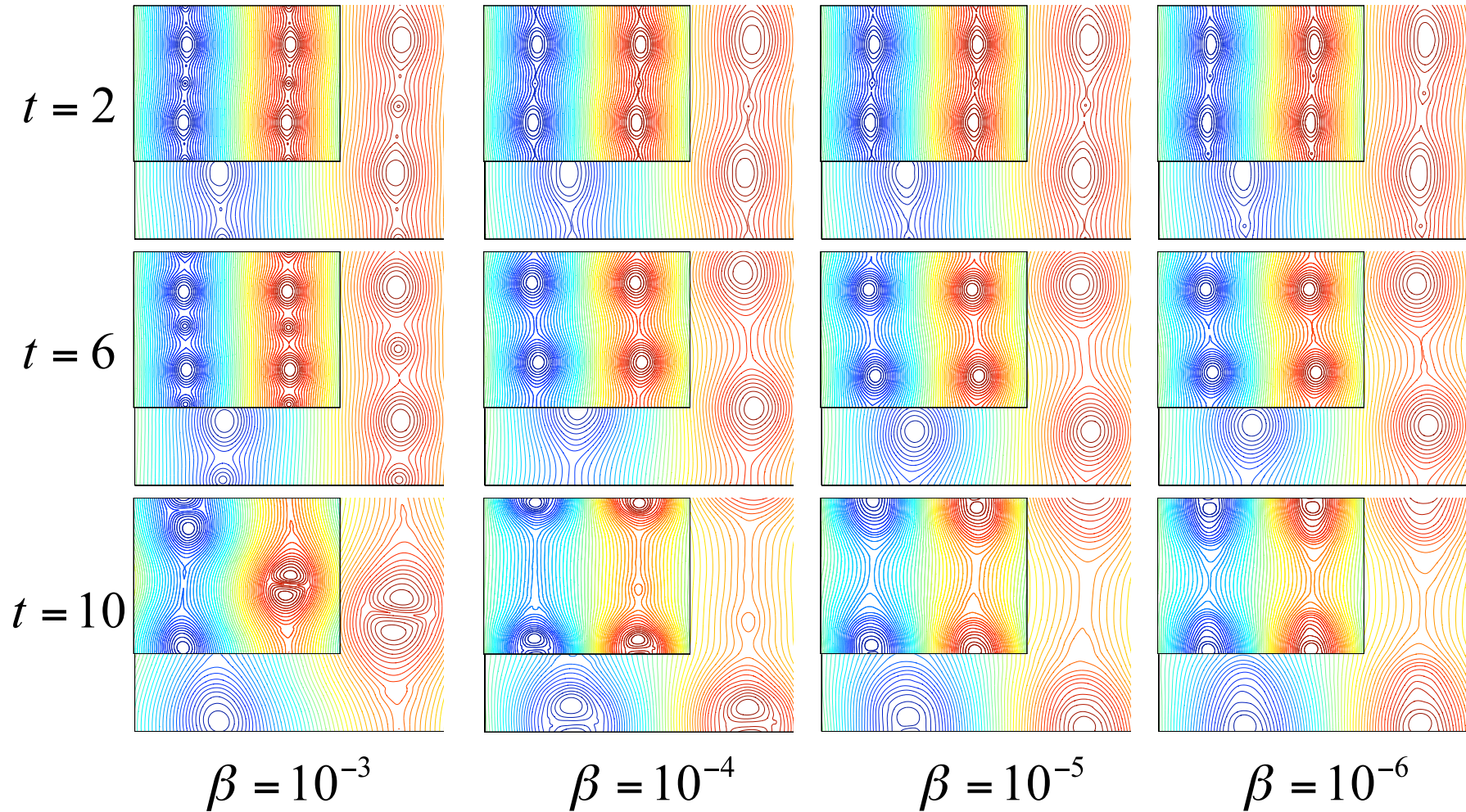




Numerical Results on Benchmarked problems (2)



Current Sheet Problem with different plasma beta values (magnetic field lines are shown)





Other Applications

