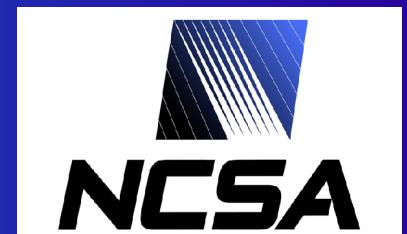


Gravity, Particles, and Cosmology with FLASH

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<http://www.astro.uiuc.edu/~pmricker/research/codes/flashcosmo/>



Overview

- The big picture
 - Applications for gravity and particles
 - Equations to be solved
- Methods
 - Algorithms currently in FLASH
 - Performance
- Usage
 - Organization of the code modules
 - Initializing a particle application
 - Analyzing particle output
 - Examples

The big picture

Gravity applications

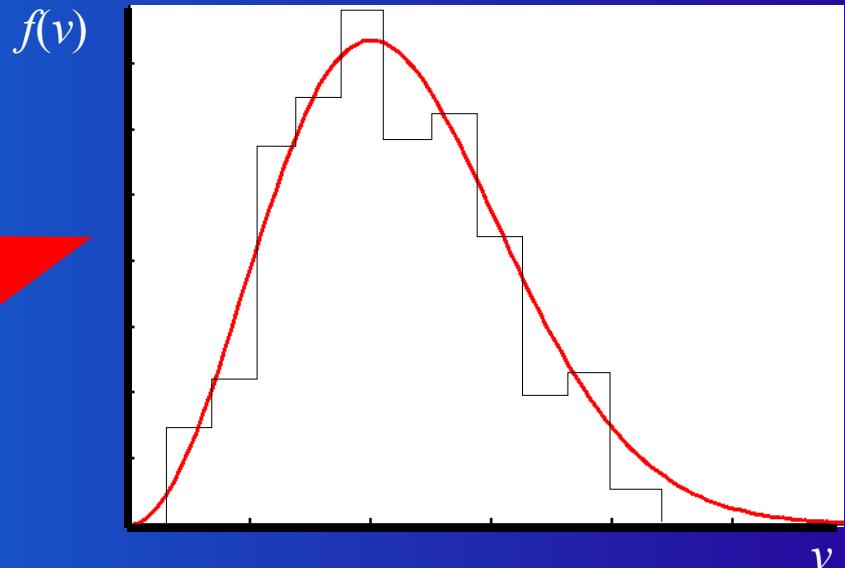
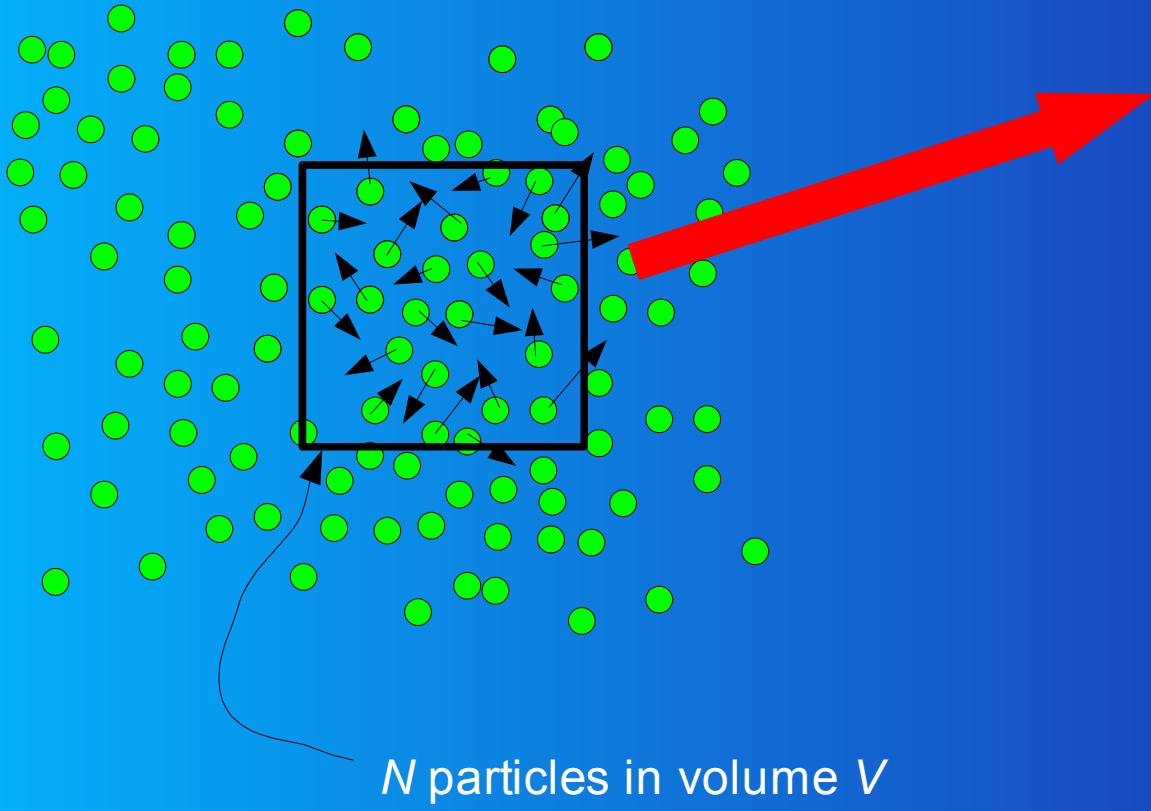
- Atmospheres (externally imposed fields)
 - Accretion (e.g., onto compact objects)
 - Buoyancy (e.g., bubbles, nuclear runaway)
- Nearly equilibrium self-gravitating systems
 - Stellar interiors
 - Intracluster medium
- Collapse and explosion problems
 - Galaxy formation
 - Star formation
 - Supernovae

Particle applications

- Dark matter
 - Large-scale structure formation
 - Galaxies
 - Galaxy clusters
- Stars
 - Galaxies
 - Globular clusters
 - Compact objects
- Tracer particles
 - Supernova nucleosynthesis
 - Mixing problems
- Electrons and ions
 - Plasmas

Particle simulation

Monte Carlo sampling of particle distribution function (gas, dust, dark matter)



$$n(\mathbf{x}) = \int f(\mathbf{x}, v) dv \approx \frac{N}{V}$$

$$\frac{|f_N - f_{\text{true}}|}{|f_{\text{true}}|} \propto N^{-1/2}$$

Basic requirements:

- As $N \rightarrow \infty$, error (“shot noise”) in approximate distribution function f_N goes to 0
- As $N \rightarrow \infty$, equation describing evolution of f_N becomes the Boltzmann equation

Equations – proper coordinates

Euler equations for gasdynamics:

Mass

$$\frac{\partial}{\partial t} \rho + \nabla \cdot (\rho \mathbf{v}) = 0$$

Momentum

$$\frac{\partial}{\partial t} (\rho \mathbf{v}) + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) + \nabla p = -\rho \nabla \phi$$

Energy

$$\frac{\partial}{\partial t} (\rho E) + \nabla \cdot [(\rho E + p) \mathbf{v}] = -\rho \mathbf{v} \cdot \nabla \phi$$

$$\rho E \equiv \frac{p}{\gamma - 1} + \frac{1}{2} \rho v^2$$

Newton's laws for particles:

$$\frac{\partial \mathbf{x}}{\partial t} = \mathbf{v}$$

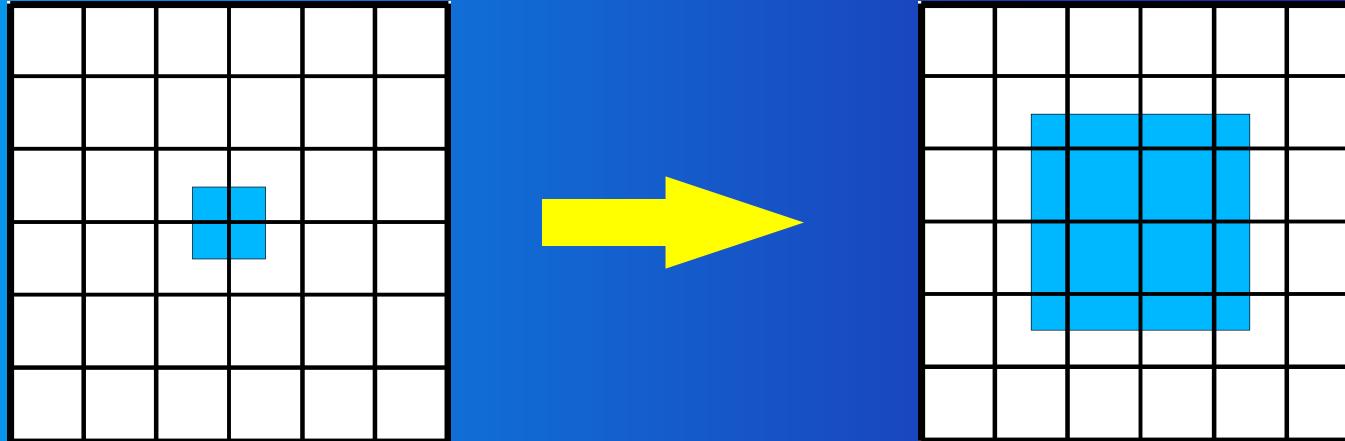
$$\frac{\partial \mathbf{v}}{\partial t} = -\nabla \phi$$

Poisson equation for gravitational potential:

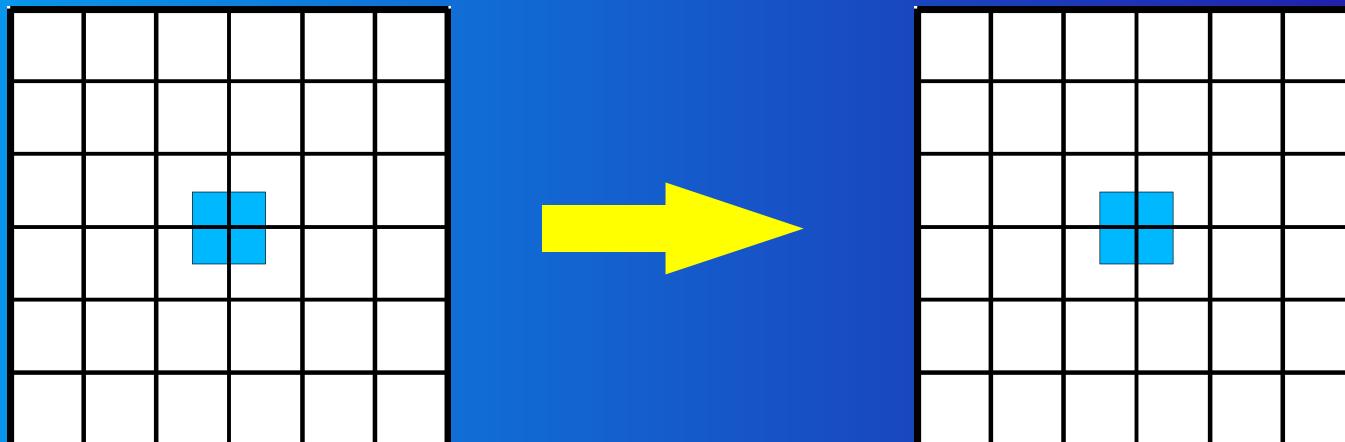
$$\nabla^2 \phi = 4\pi G \rho$$

Comoving coordinates

Proper coordinates \mathbf{r} are Eulerian: $\mathbf{u} = \dot{\mathbf{r}}$



Comoving coordinates \mathbf{x} scale out the expansion: $\mathbf{v} = \dot{\mathbf{x}}$



$$\mathbf{r} = a \mathbf{x} \Rightarrow \mathbf{u} = a \dot{\mathbf{x}} + \dot{a} \mathbf{x} = a \mathbf{v} + H \mathbf{r} \quad H \equiv \frac{\dot{a}}{a} \text{ is the Hubble parameter}$$

Fluid equations – comoving coordinates

Comoving variables: $\rho \equiv a^3 \tilde{\rho}$, $p \equiv a \tilde{p}$, $T \equiv \frac{\tilde{T}}{a^2}$

Fluid equations for gasdynamics:

Mass

$$\frac{\partial}{\partial t} \rho + \nabla \cdot (\rho \mathbf{v}) = 0$$

Momentum $\frac{\partial}{\partial t} (\rho \mathbf{v}_g) + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) + \nabla p + 2 \frac{\dot{a}}{a} \rho \mathbf{v} = -\rho \nabla \phi$

Energy $\frac{\partial}{\partial t} (\rho E) + \nabla \cdot [(\rho E + p) \mathbf{v}] + \frac{\dot{a}}{a} \left[\left(\frac{3\gamma-1}{\gamma-1} \right) p + 2 \rho v^2 \right] = \rho \frac{\partial \phi}{\partial t}$

$$\rho E \equiv \frac{p}{\gamma-1} + \frac{1}{2} \rho v^2 + \rho \phi$$

Particle/field equations – comoving

Particle equations:

$$\frac{\partial \mathbf{x}}{\partial t} = \mathbf{v}$$

$$\frac{\partial \mathbf{v}}{\partial t} + 2 \frac{\dot{a}}{a} \mathbf{v} = -\nabla \phi$$

Gravity (Poisson & Friedmann equations):

$$\nabla^2 \phi = \frac{4\pi G}{a^3} [(\rho_g + \rho_{dm}) - \overline{(\rho_g + \rho_{dm})}]$$

$$H^2(t) \equiv \left(\frac{\dot{a}}{a} \right)^2 = H_0^2 \left(\frac{\Omega_m}{a^3} + \frac{\Omega_r}{a^4} + \Omega_\Lambda - \frac{\Omega_c}{a^2} \right)$$

Methods

Gravity algorithms in FLASH

- Externally imposed fields
 - Constant field ($\mathbf{g} = \text{constant}$)
 - Point mass ($\mathbf{g} = -GM\mathbf{r}/r^3$)
 - Plane-parallel ($\mathbf{g} = -GM\mathbf{x}/r^3$)
- Self-gravity – Poisson equation
 - Multipole
 - Nearly spherically symmetric problems
 - Isolated boundary conditions
 - 1D spherical, 2D cylindrical, 3D Cartesian mesh geometries
 - Multigrid
 - Arbitrary source fields
 - Periodic or isolated boundary conditions (James algorithm)
 - 1D/2D/3D Cartesian mesh geometry

Multipole algorithm

Solve Poisson's equation in integral form:

$$\phi(\mathbf{x}) = -G \iiint d^3x' \frac{\rho(\mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} - \frac{G}{|\mathbf{x} - \mathbf{x}'|} = \text{Green's function}$$

Note that

$$\frac{1}{|\mathbf{x} - \mathbf{x}'|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^{l} \frac{1}{2l+1} \frac{r_<^l}{r_>} Y_{lm}^*(\theta', \varphi') Y_{lm}(\theta, \varphi)$$

where the Y_{lm} are spherical harmonic basis functions and

$$r_< \equiv \min[|\mathbf{x}|, |\mathbf{x}'|]$$
$$r_> \equiv \max[|\mathbf{x}|, |\mathbf{x}'|]$$

Thus

$$\phi(\mathbf{x}) = -G \sum_{l=0}^{\infty} \sum_{m=-l}^{l} \frac{1}{2l+1} Y_{lm}(\theta, \varphi) \left[r^l \int_{r < r'} d^3x' \frac{\rho(\mathbf{x}')}{r'^{l+1}} Y_{lm}^*(\theta', \varphi') + \frac{1}{r^{l+1}} \int_{r > r'} d^3x' \rho(\mathbf{x}') r'^l Y_{lm}^*(\theta', \varphi') \right]$$

Multipole algorithm

Using definition of Y_{lm} in terms of Legendre polynomials P_{lm} , obtain

$$\begin{aligned}\phi(\mathbf{x}) = & -G \sum_{l=0}^{\infty} P_{l0}(\cos \theta) \left[r^l \mu_{l0}^{eo}(r) + \frac{1}{r^{l+1}} \mu_{l0}^{ei}(r) \right] - \\ & 2G \sum_{l=1}^{\infty} \sum_{m=1}^l P_{lm}(\cos \theta) \left[(r^l \cos m\varphi) \mu_{lm}^{eo}(r) + (r^l \sin m\varphi) \mu_{lm}^{oo}(r) + \right. \\ & \left. \frac{\cos m\varphi}{r^{l+1}} \mu_{lm}^{ei}(r) + \frac{\sin m\varphi}{r^{l+1}} \mu_{lm}^{oi}(r) \right]\end{aligned}$$

where the even/odd, inner/outer moments of the density are given by

$$\mu_{lm}^{ei}(r) \equiv \frac{(l-m)!}{(l+m)!} \int_{r>r'} d^3x' r'^l \rho(\mathbf{x}') P_{lm}(\cos \theta') \cos m\varphi'$$

$$\mu_{lm}^{oi}(r) \equiv \frac{(l-m)!}{(l+m)!} \int_{r>r'} d^3x' r'^l \rho(\mathbf{x}') P_{lm}(\cos \theta') \sin m\varphi'$$

$$\mu_{lm}^{eo}(r) \equiv \frac{(l-m)!}{(l+m)!} \int_{r<r'} d^3x' r'^{-(l+1)} \rho(\mathbf{x}') P_{lm}(\cos \theta') \cos m\varphi'$$

$$\mu_{lm}^{oo}(r) \equiv \frac{(l-m)!}{(l+m)!} \int_{r<r'} d^3x' r'^{-(l+1)} \rho(\mathbf{x}') P_{lm}(\cos \theta') \sin m\varphi'$$

These moments can be evaluated by any appropriate $O(N)$ quadrature method.

Relaxation methods

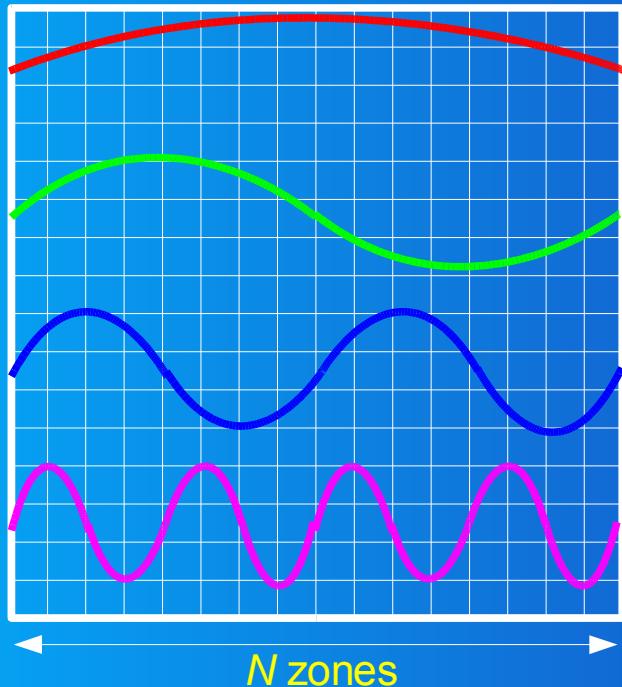
Treat differential form of Poisson equation as steady-state solution to a diffusion problem:

$$\nabla^2 \phi = \rho \quad \rightarrow \quad \frac{\partial \phi}{\partial t} = \nabla^2 \phi - \rho$$

“Time coordinate” is really an iteration coordinate. Difference explicitly, rearrange, and apply stability criterion to get Jacobi method:

$$\phi_i^{n+1} = \frac{1}{2} [\phi_{i+1}^n + \phi_{i-1}^n] - \frac{\Delta x^2}{2} \rho_i$$

Simple... but $O(N^2)$!



2N iterations to communicate across wavelength

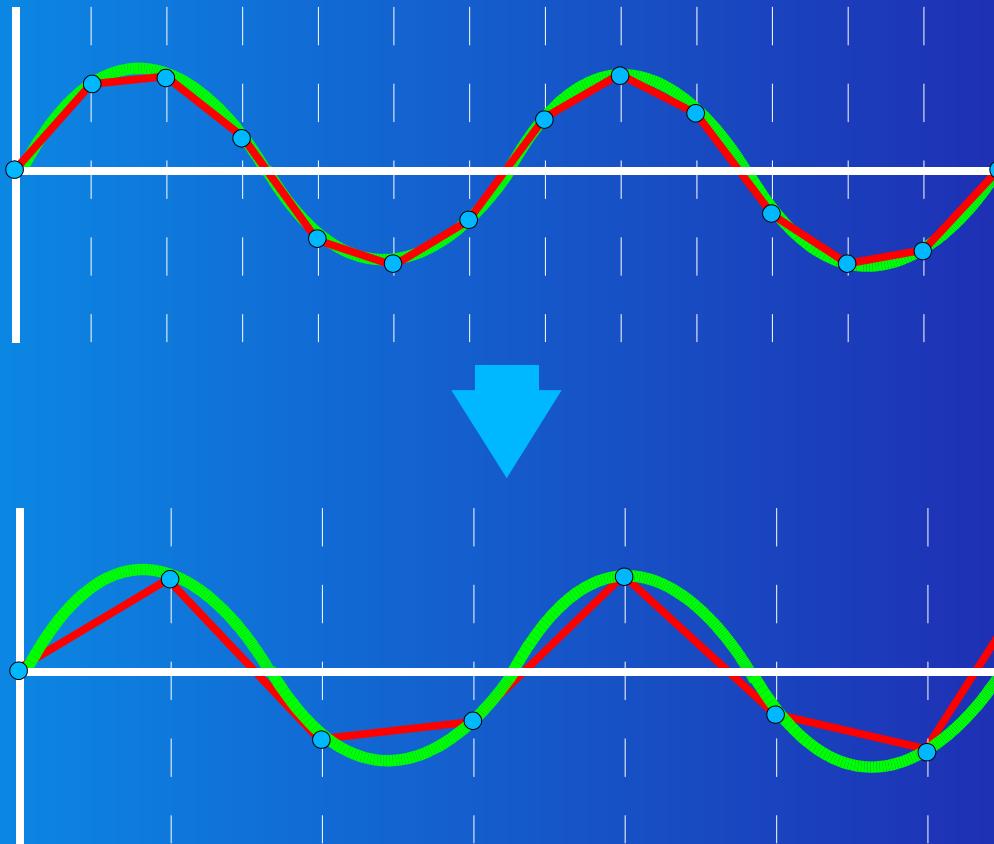
N iterations to communicate across wavelength

$N/2$ iterations to communicate across wavelength

$N/4$ iterations to communicate across wavelength

Multigrid – basic idea

- On a coarser mesh, a given error mode appears to be “higher frequency:”



- A single multigrid iteration (“V cycle” in our case) comprises
 - Relaxation of long-wavelength error modes on coarse mesh
 - Relaxation of short-wavelength error modes on fine mesh
 - Both wavelengths converge at the same rates – so $O(N \log N)$ performance

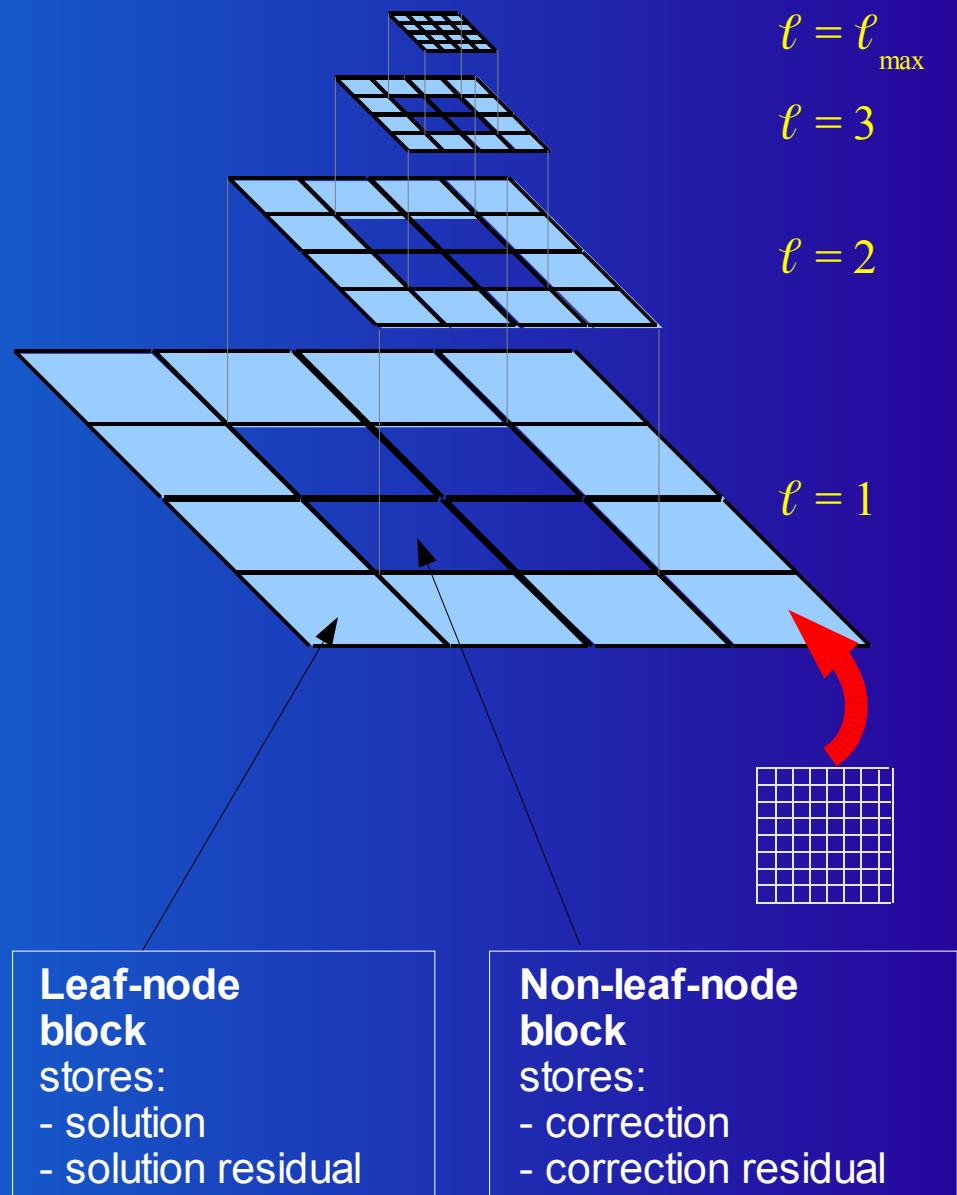
AMR multigrid – Poisson equation

Based on Martin & Cartwright (1996)

1. Compute residual of the source equation on all leaf blocks. **Residual operation uses flux conservation at fine-coarse boundaries.**

2. V-cycle:

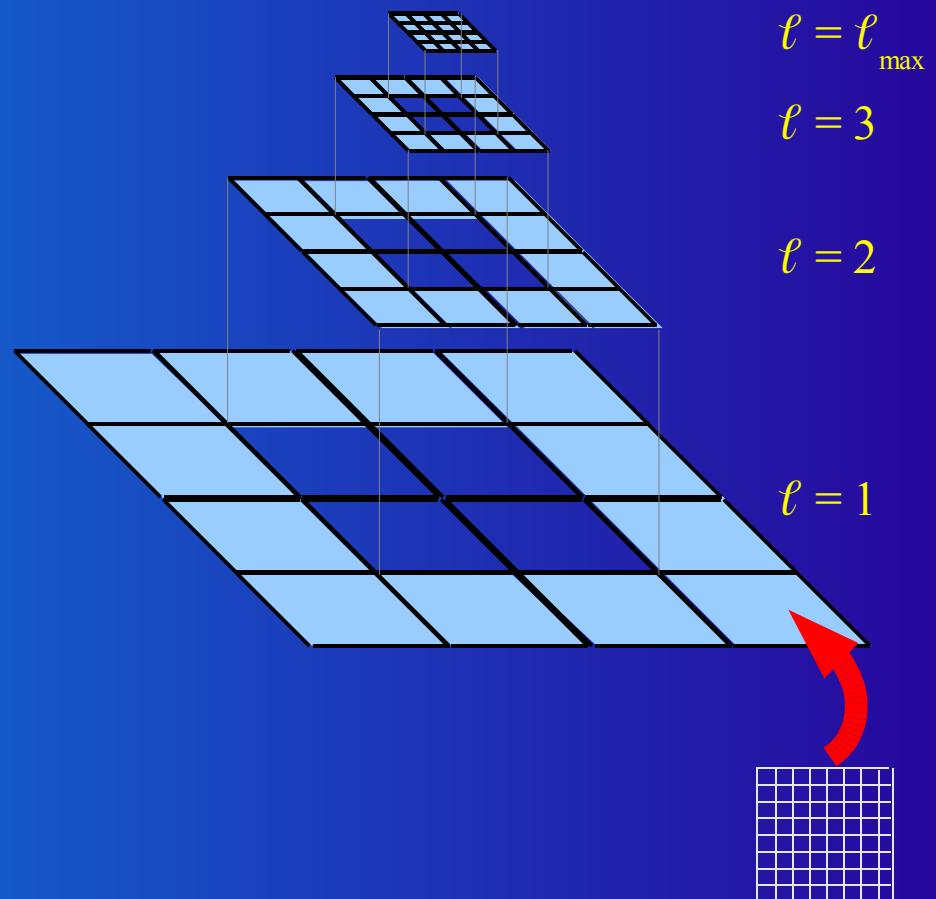
- a) Zero the correction on level ℓ_{\max} .
- b) For $\ell = \ell_{\max}$ down to 2:
 - i. Copy solution to a temporary variable.
 - ii. Zero correction on level $\ell-1$.
 - iii. Relax correction equation on level ℓ .
 - iv. Add correction to the solution on level ℓ .
 - v. Compute residual of the correction equation on all blocks (leaf or not) on level ℓ . Restrict this residual to level $\ell-1$.
 - vi. Compute the residual of the source equation on all leaf blocks of level $\ell-1$.
- c) Solve correction equation on level 1.
Correct the solution on this level.



AMR multigrid – Poisson equation

2. d) For $\ell = 2$ up to ℓ_{\max} :

- i. Prolongate correction from level $\ell-1$ and add result to the correction on level ℓ .
- ii. Replace the stored residual on level ℓ with the new residual of the correction equation.
- iii. Zero a second temporary variable on levels $\ell-1$ and ℓ .
- iv. Relax this variable against the residual on level ℓ .
- v. Add the result to the correction on level ℓ .
- vi. Copy back into the solution on all leaf blocks on level ℓ .
- vii. Add the correction to the solution on all leaf blocks on level ℓ .



Important to use quadratic ($O(\Delta x^3)$) interpolation when prolongating, or V-cycles will fail to converge (interpolation error will be comparable to differencing error).

Pieces of a particle simulation algorithm

Given particle positions \mathbf{x}_i^n and velocities \mathbf{v}_i^n at time t_n , compute values at time t_{n+1} :

1. Compute acceleration of each particle at t_n : \mathbf{a}_i^n
2. Using \mathbf{a}_i^n and possibly acceleration values from previous steps, advance \mathbf{v}_i^n :

$$\mathbf{v}_i^{n+1} = \mathbf{v}_i^n + \Delta t \mathbf{a}_{\text{eff},i}(\mathbf{a}_i^n, \dots)$$

3. Using \mathbf{v}_i^n and possibly velocity values from previous steps, advance \mathbf{x}_i^n :

$$\mathbf{x}_i^{n+1} = \mathbf{x}_i^n + \Delta t \mathbf{v}_{\text{eff},i}(\mathbf{v}_i^n, \dots)$$

Positions and velocities may be evaluated at different times (e.g., t_n and $t_{n+1/2}$).

Basic components:

- Field computation
 - Direct N -body (particle-particle)
 - Particle-mesh
 - Particle-particle-particle mesh (P^3M)
 - Treecodes
- Time integrator

Particle algorithms in FLASH

- Passive particles – particles move with gas
 - Time integrators
 - Euler
 - Predictor-corrector
- Active particles – move independently of gas
 - Time integrators
 - Euler
 - Leapfrog
 - Cosmological leapfrog
 - Long-range forces
 - Gravity (particle-mesh)
- Particle-mesh transfer operators
 - Nearest grid point (NGP)
 - Cloud in cell (CIC)
 - Triangle-shaped cloud (TSC)

Particle-mesh method

- Exploit fast mesh-based Poisson solvers
- First introduced for plasma simulation in the early 1960s
- Procedure:
 1. Assign particle masses to mesh using a mass assignment operator $\rightarrow \rho_{ijk}$.
 2. Solve $\nabla^2\phi = 4\pi G\rho$ on mesh.
 3. Finite-difference ϕ to get forces on mesh.
 4. Interpolate forces to the particle positions using a force interpolation operator.
 5. Advance the positions and velocities of the particles in time.
 6. Repeat.

Particle-mesh transfer operators

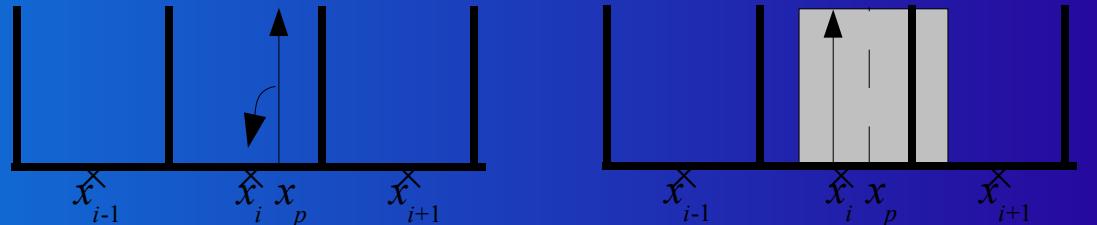
Mass assignment operator:

$$\rho_{ijk} = \frac{m}{\Delta x \Delta y \Delta z} \sum_{p=1}^{N_p} W(x_i - x_p) W(y_j - y_p) W(z_k - z_p)$$

Force interpolation operator:

$$F(x_p) = -m \sum_{i=1}^{N_g} \sum_{j=1}^{N_g} \sum_{k=1}^{N_g} (\nabla \phi)_{ijk} W(x_p - x_i) W(y_p - y_j) W(z_p - z_k)$$

Nearest grid point (NGP)



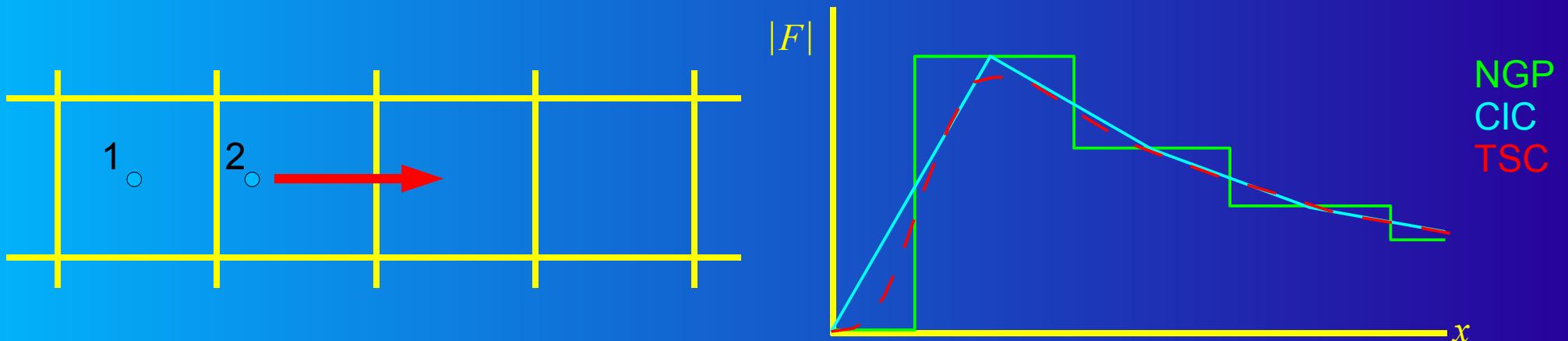
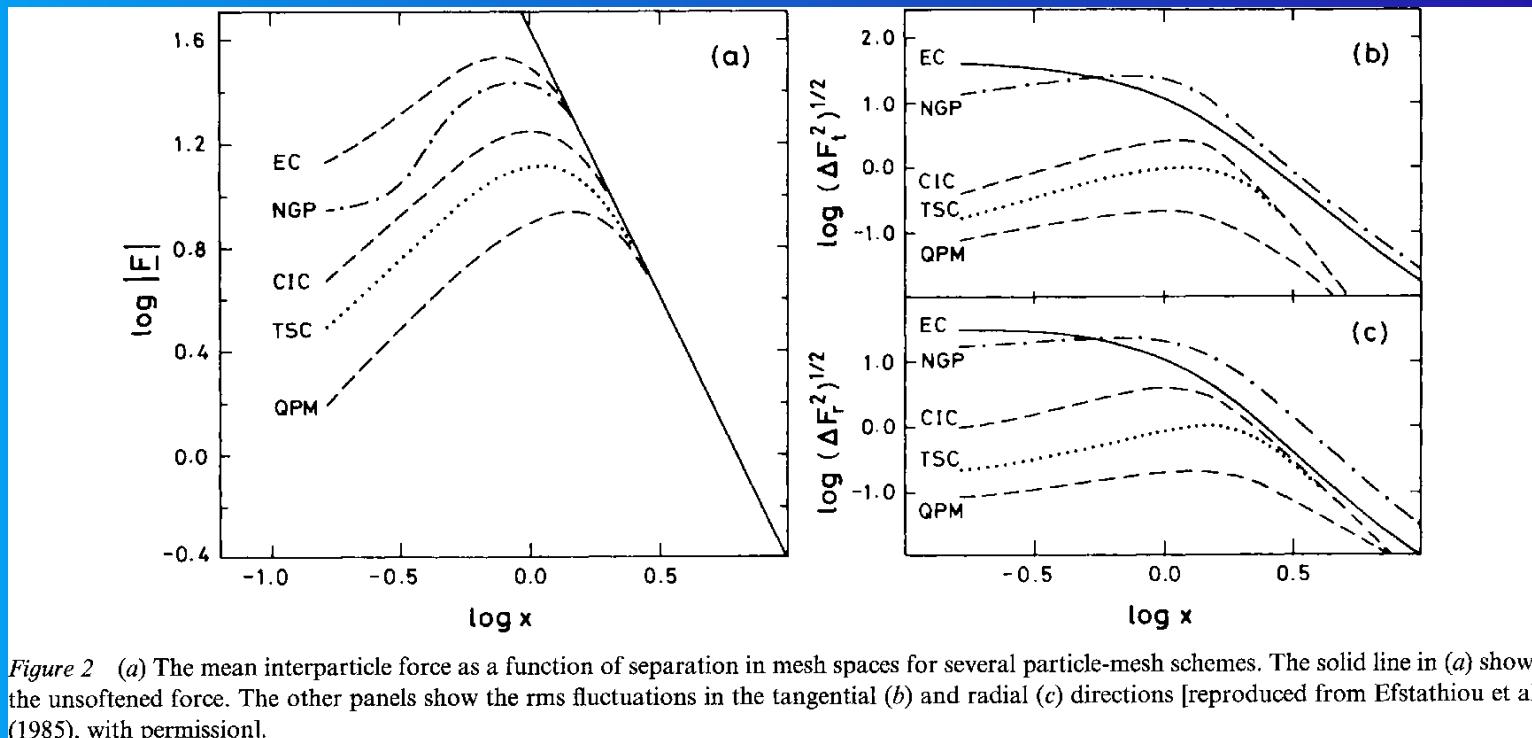
Cloud in cell (CIC)



Triangle-shaped cloud (TSC)



Force smoothing in particle-mesh



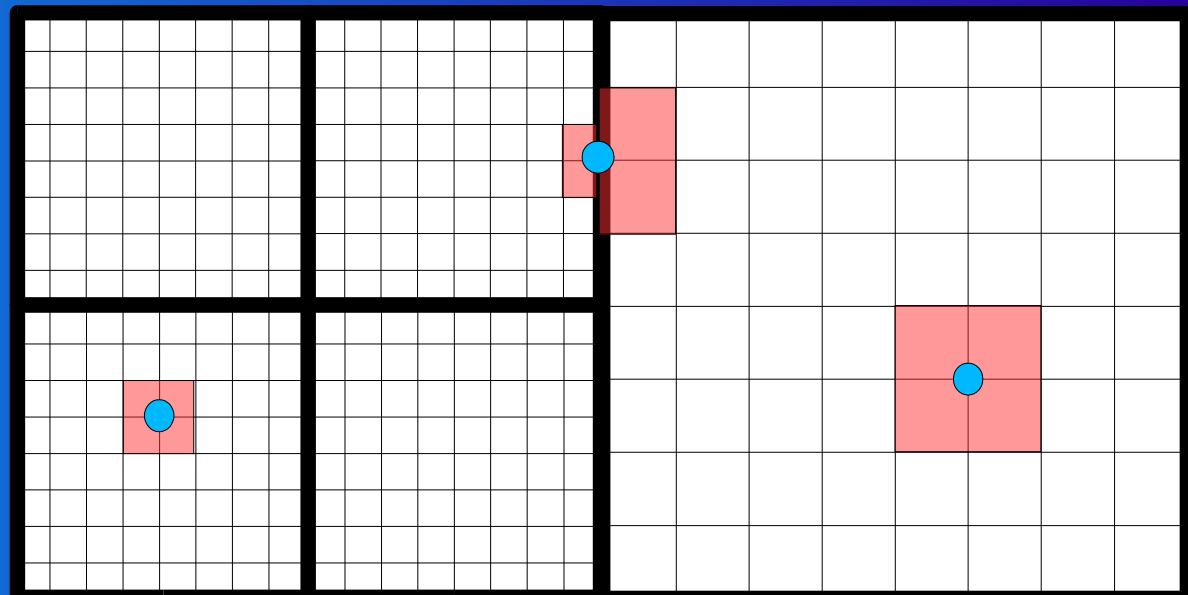
At jumps in refinement

- Multigrid solver
- Quadratic interpolants
- “Flux” matching
- Particle interpolation
- Particle clouds change discontinuously at boundary
- Nonzero self-forces
- Few particles: refine on mesh particle density
- Many particles: mean field dominates

$$\frac{\partial^2 \phi}{\partial x^2} = \rho \rightarrow g_{i+1/2,j} = \frac{\phi_{i+1,j} - \phi_{i,j}}{\Delta x}$$

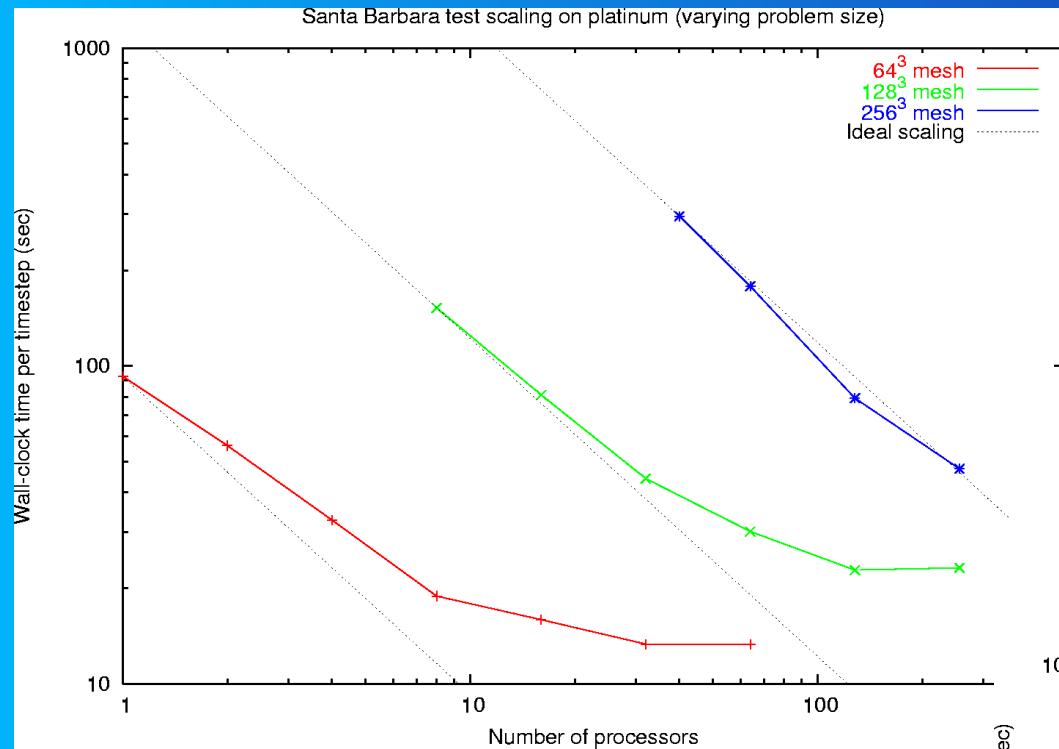
$$g_{1/2,j}^{\text{coarse}} := g_{N+1/2,j}^{\text{fine}}$$

$$R_{i,j} = \rho_{i,j} - \frac{g_{i+1/2,j} - g_{i-1/2,j}}{\Delta x}$$

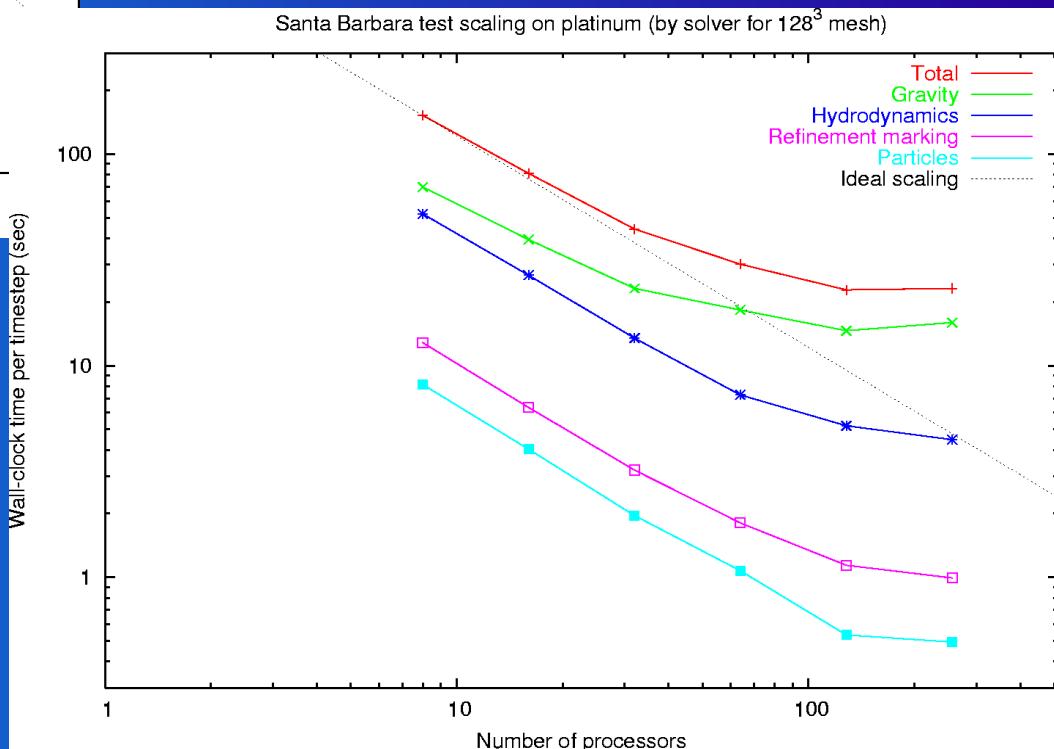


Scaling by solver

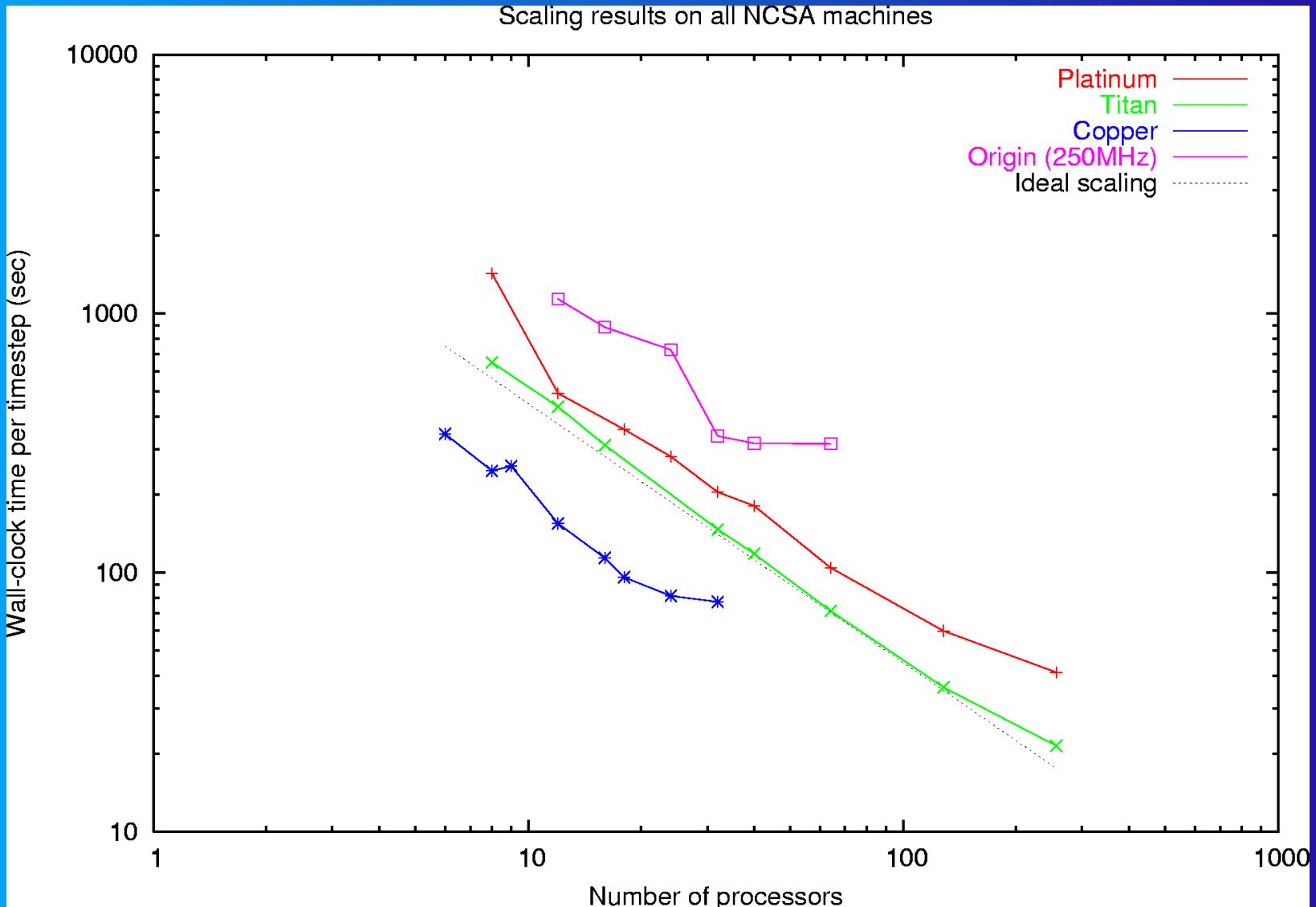
Uniform mesh on NCSA Platinum IA-32 cluster (1 GHz PIII CPUs)



- 16^3 zones per block
- No effort at storage balancing (uniform block weights)

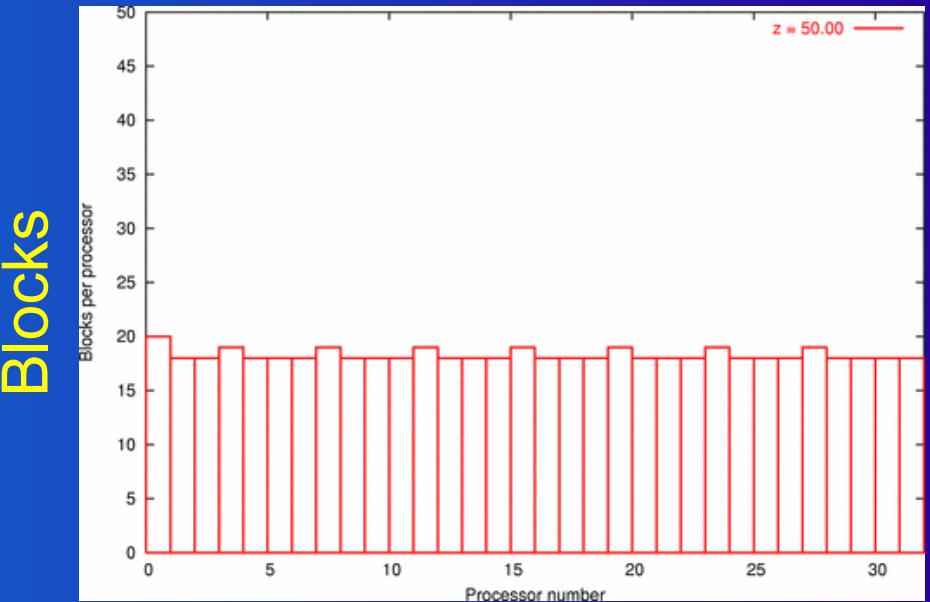


Scaling by platform

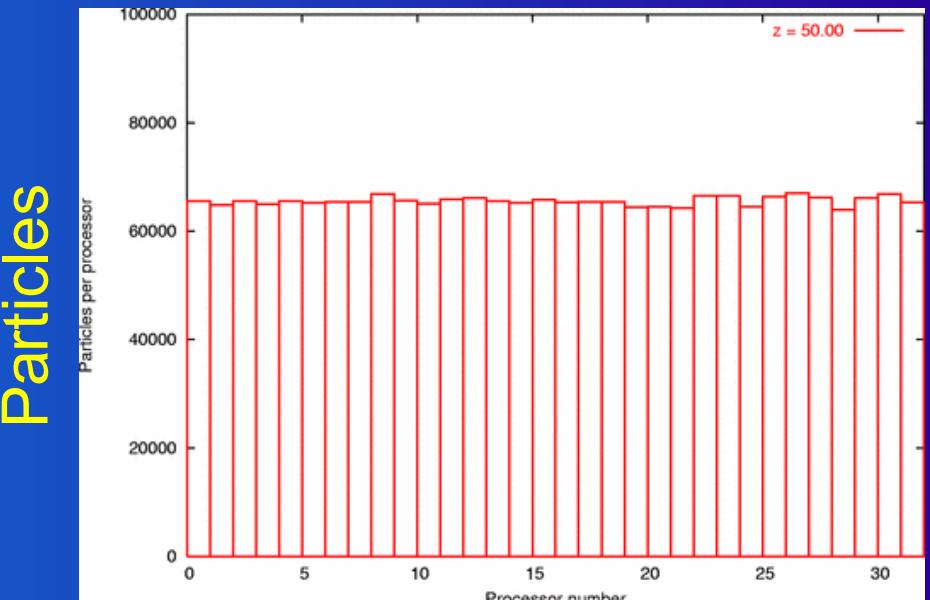


Storage balancing

- The problem
 - Particles start uniform, end clustered
 - Locality: particles stored with blocks
 - Uniform block weights distribute blocks well, but not particles



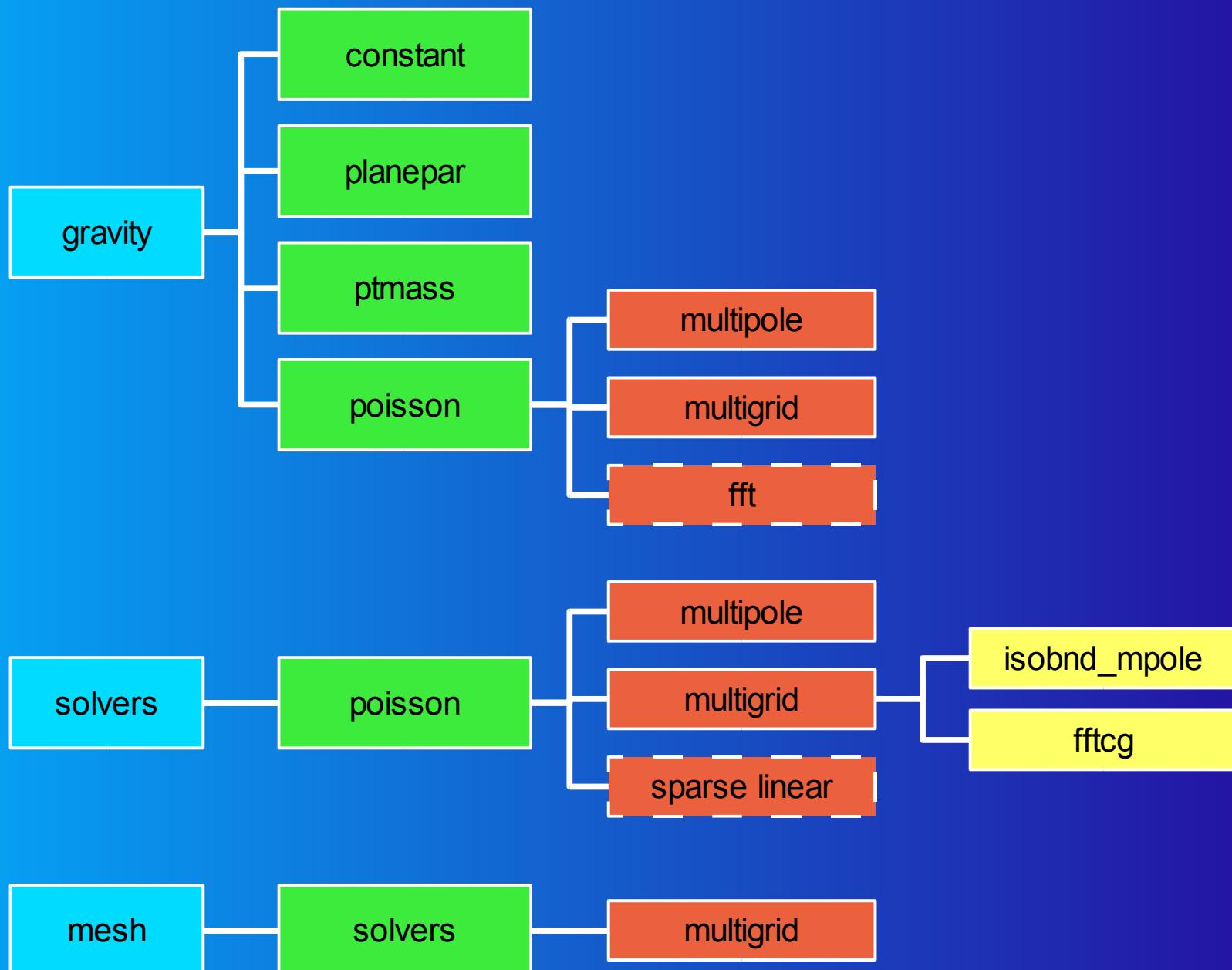
- A solution
 - Weight blocks by number of particles they contain
 - Shifts storage imbalance to blocks
 - Density-based refinement optimizes both distributions



Processor number

Usage

FLASH gravity-related modules



Methods supplied by Gravity F90 module

- Usage:

```
use Gravity, ONLY: <list of gravity methods to import>
```

- Types of sub-modules in FLASH

- Compute acceleration “natively”
- Compute potential, then difference to get acceleration

- Methods generic to all gravity sub-modules (not called by user modules)

- `InitGravity()`

Initialize the gravity module (called by `init_flash()`)

- `GravityTimestep(dt_grav, dt_minloc, block)`

Compute any timestep restrictions imposed by gravity module (called by `timestep()`)

Methods supplied by Gravity F90 module

- Methods oriented toward computing gravitational acceleration
- No mesh variables generically defined; declared by physics FLASH modules
 - `GravAccelAllBlocks(pot_var, acc_var, component)`
Compute a component of the acceleration on all blocks.
 - `GravAccelOneBlock(pot_var, acc_var, component, block)`
Compute a component of the acceleration on a single block.
 - `GravAccelOneLevel(pot_var, acc_var, component, level)`
Compute a component of the acceleration on all blocks at a single level of refinement.
 - `GravAccelOneRow(j, k, direction, block, pot_var, g, n)`
Compute a component of the acceleration along a single row of a single block; return as an array (`g`).
 - `GravAccelOneZone(x, y, z, g)`
Compute all components of the acceleration at a single position in space.

Methods supplied by Gravity F90 module

- Methods oriented toward computing gravitational potential
- Generically defined mesh variables: gpot, gpol
 - GravPotentialAllBlocks (pot_var)
Compute potential on the entire mesh.
 - GravPotentialOneBlock (pot_var, block)
(Planned) return potential on a single block.
 - GravPotentialOneLevel (pot_var, level)
(Planned) return potential on all blocks at a given level of refinement.

Gravity FLASH module configuration

```
# Parameters:

D      grav_boundary          External boundary condition to use for Poisson solver
D          &                      (applied to all boundaries): 0 = isolated,
D          &                      1 = periodic, 2 = Dirichlet
D      grav_boundary_type      String-valued version of grav_boundary. Accepts:
D          &                      "isolated", "periodic", "dirichlet".
D          &                      If grav_boundary is set, its value overrides this
D          &                      setting.
D      igrav                  Gravity switch: if /= 0, turn on gravity

DEFAULT multigrid

EXCLUSIVE multigrid fft multipole

PARAMETER grav_boundary_type STRING "isolated" # string-valued boundary parm
PARAMETER grav_boundary      INTEGER -1        # integer-valued boundary parm
PARAMETER igrav              INTEGER 1        # if /= 0, turn on gravity

# Self-gravity requires that the density be defined as a variable.
# This should be ignored by setup if the hydro module (which also
# defines "dens") is included.

VARIABLE gpot NOADVECT NORENORM NOCONSERVE # grav. potential at current step
VARIABLE gpol NOADVECT NORENORM NOCONSERVE # grav. potential at previous step
VARIABLE dens ADVECT NORENORM CONSERVE    # density
```

Multigrid FLASH module configurations

solvers/poisson/multigrid/isobnd_mpole (isolated boundaries via James method)

D mpole_lmax Maximum multipole moment to use

PARAMETER mpole_lmax INTEGER 0

mesh/solvers/multigrid

D mgrid_max_residual_norm Maximum ratio of the residual norm to that of the right-hand side
D mgrid_max_iter_change Maximum change in the residual norm from one iteration to the next
D mgrid_max_vcycles Maximum number of V-cycles to take
D mgrid_npresmooth Number of pre-smoothing iterations to perform on each level
D mgrid_npostsmooth Number of post-smoothing iterations to perform on each level
D mgrid_smooth_tol Convergence criterion (coarse grid smoother)
D mgrid_solve_max_iter Maximum number of iterations for solution (coarse grid smoother)
D mgrid_print_norm If .true., print residual norm to stdout after each V-cycle
D quadrant In 2d cylindrical coords, assume symmetry about y=0

PARAMETER mgrid_max_residual_norm REAL 1.E-6

PARAMETER mgrid_max_iter_change REAL 1.E-3

PARAMETER mgrid_max_vcycles INTEGER 100

PARAMETER mgrid_npresmooth INTEGER 1

PARAMETER mgrid_npostsmooth INTEGER 8

PARAMETER mgrid_smooth_tol REAL 5.E-3

PARAMETER mgrid_solve_max_iter INTEGER 5000

PARAMETER mgrid_print_norm BOOLEAN FALSE

PARAMETER quadrant BOOLEAN FALSE

Work variables needed by multigrid solver

VARIABLE mgw1 NOADVECT NORENORM NOCONSERVE

VARIABLE mgw2 NOADVECT NORENORM NOCONSERVE

VARIABLE mgw3 NOADVECT NORENORM NOCONSERVE

VARIABLE mgw4 NOADVECT NORENORM NOCONSERVE

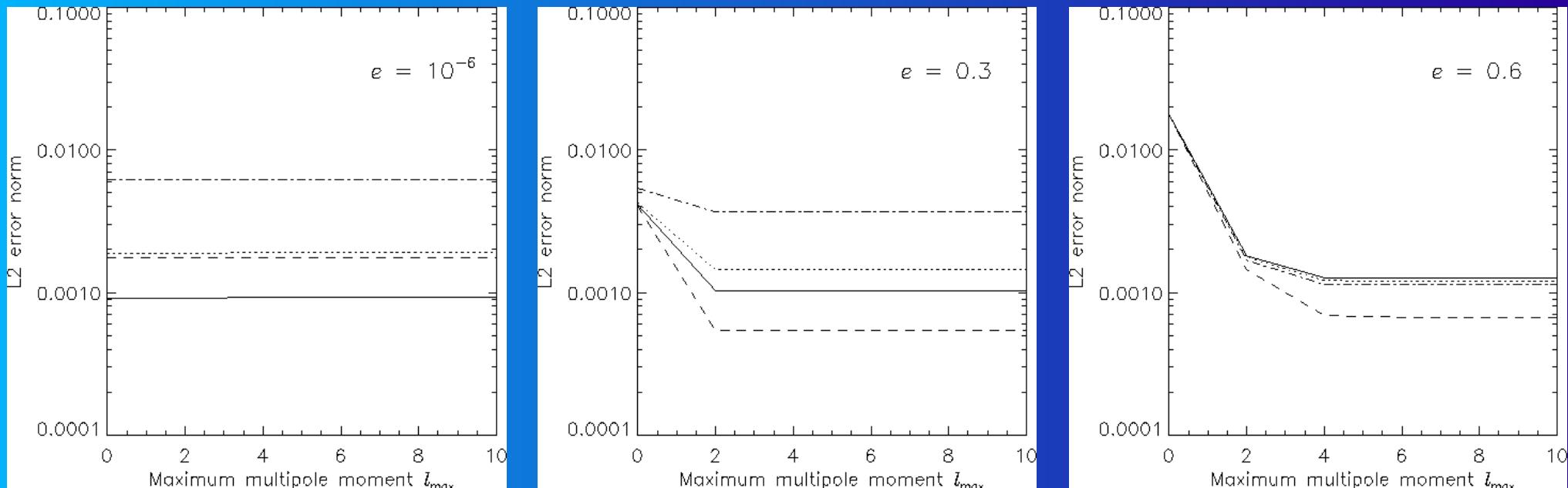
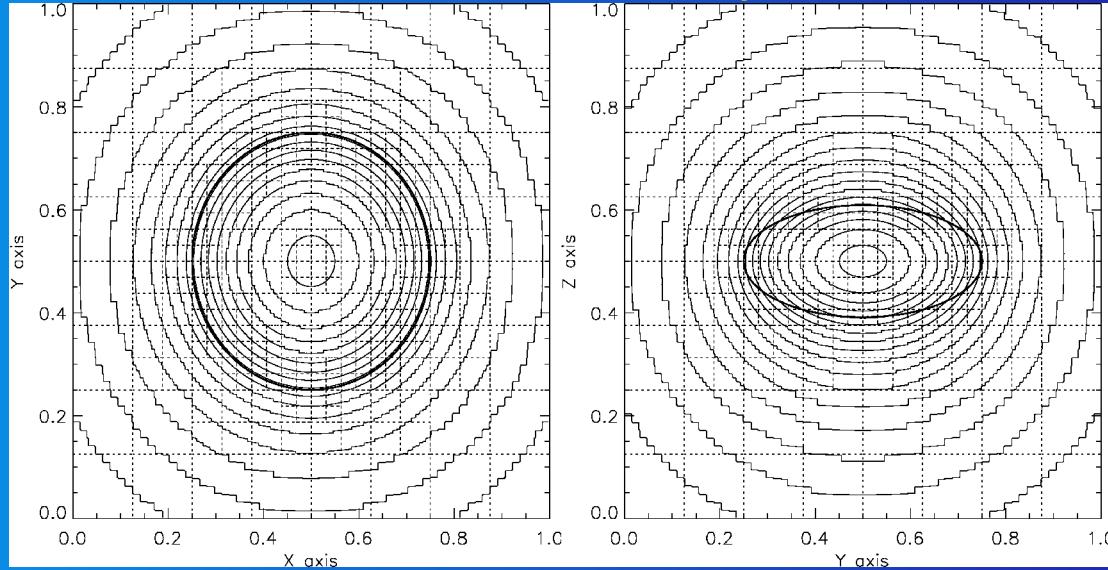
VARIABLE mgw5 NOADVECT NORENORM NOCONSERVE

VARIABLE mgw6 NOADVECT NORENORM NOCONSERVE

VARIABLE mgw7 NOADVECT NORENORM NOCONSERVE

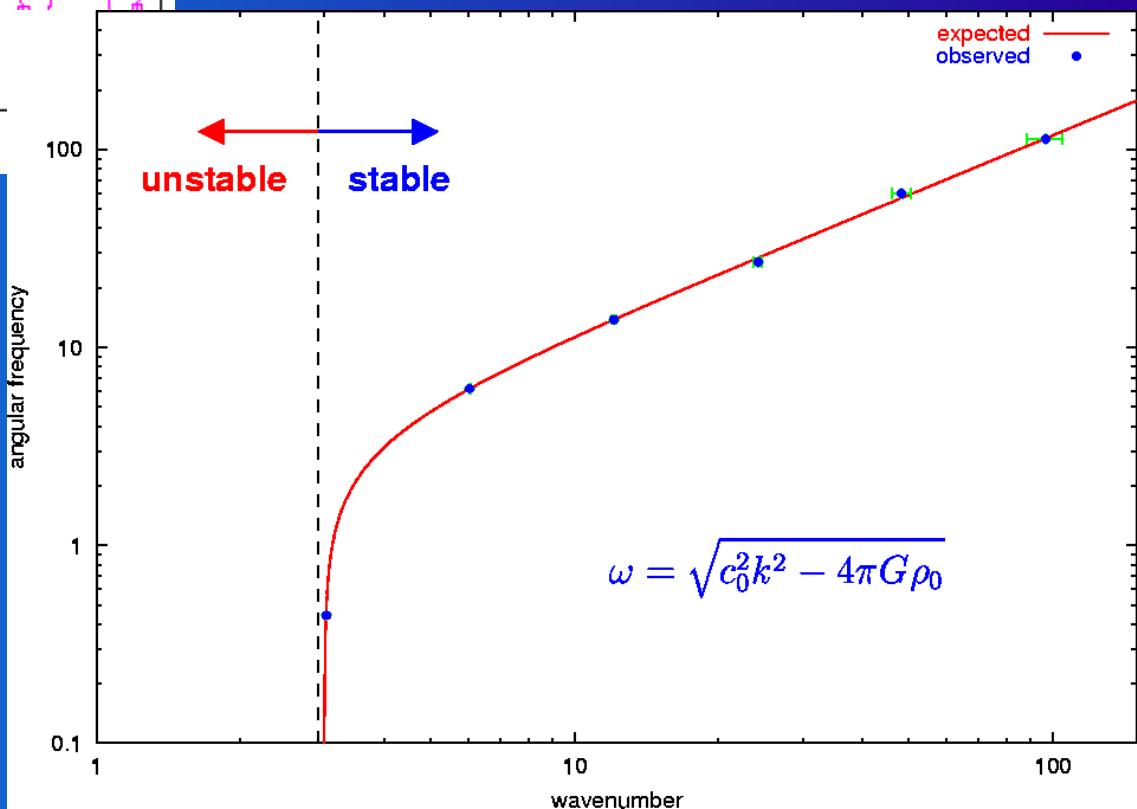
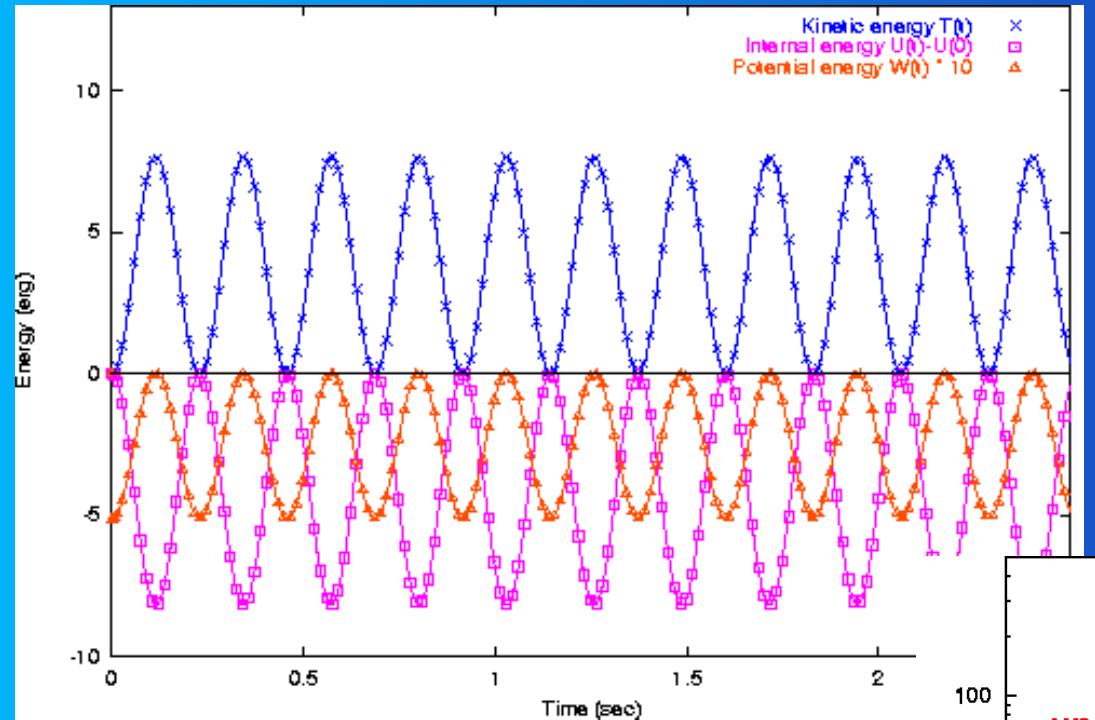
Example: Maclaurin spheroid potential

Eccentricity 0.9, $\ell_{\max} = 10$

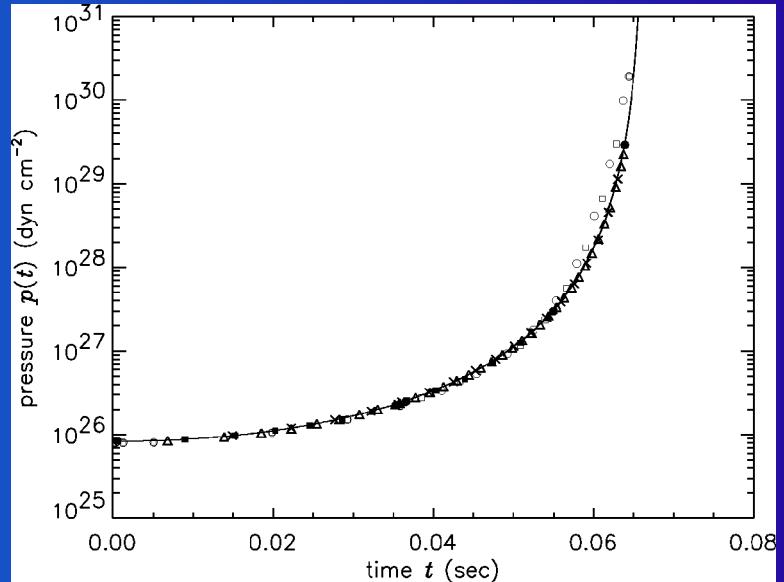
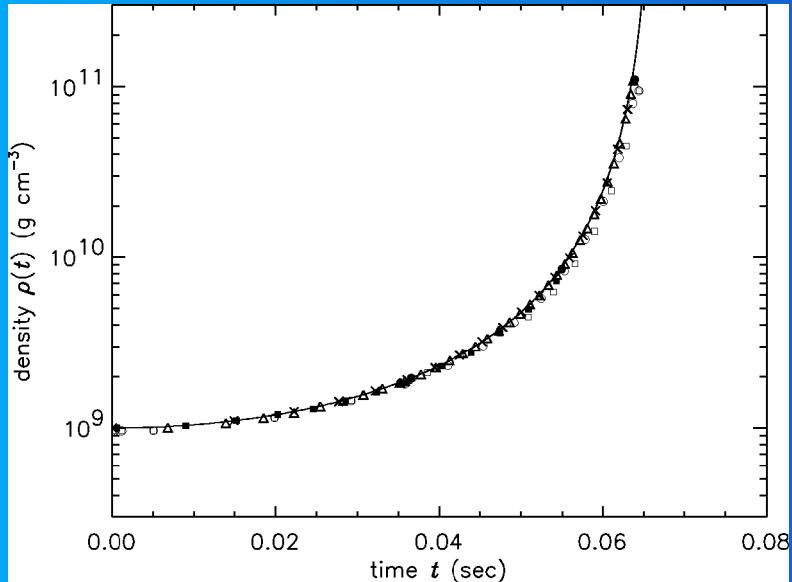
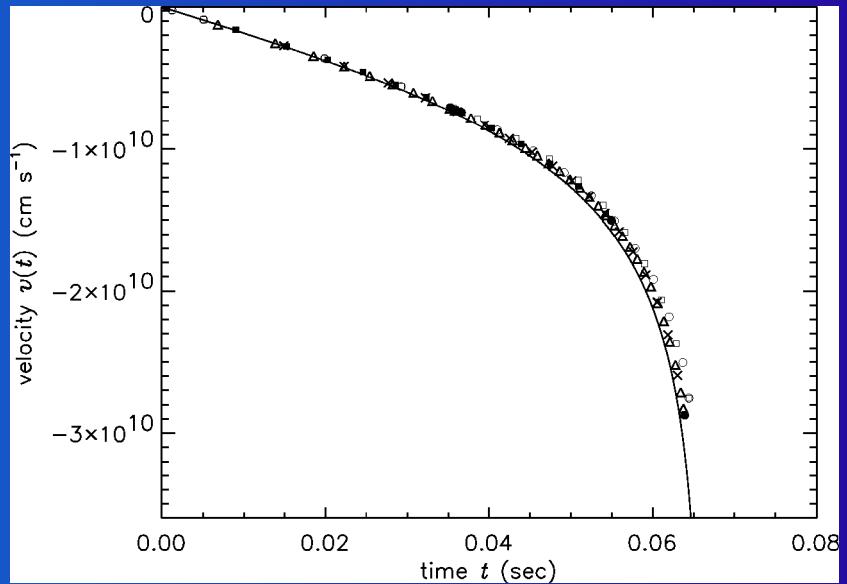
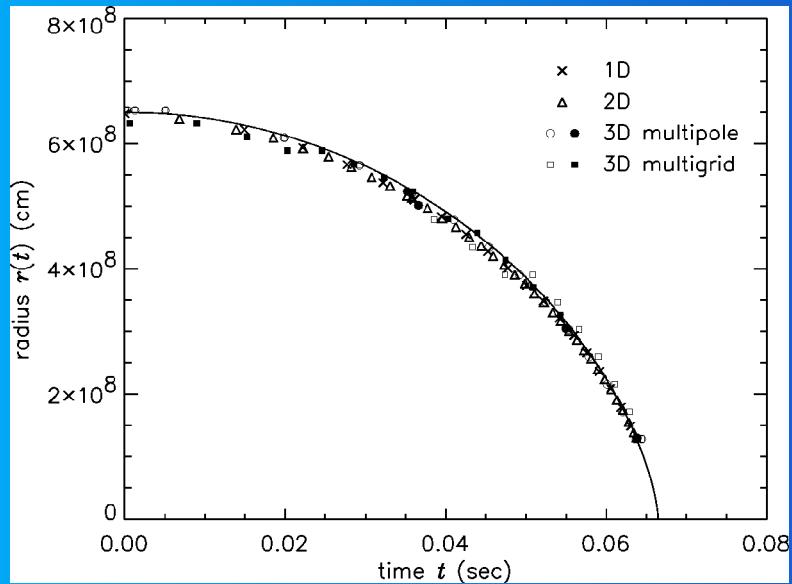


Varying eccentricity and multipole order

Example: Jeans instability

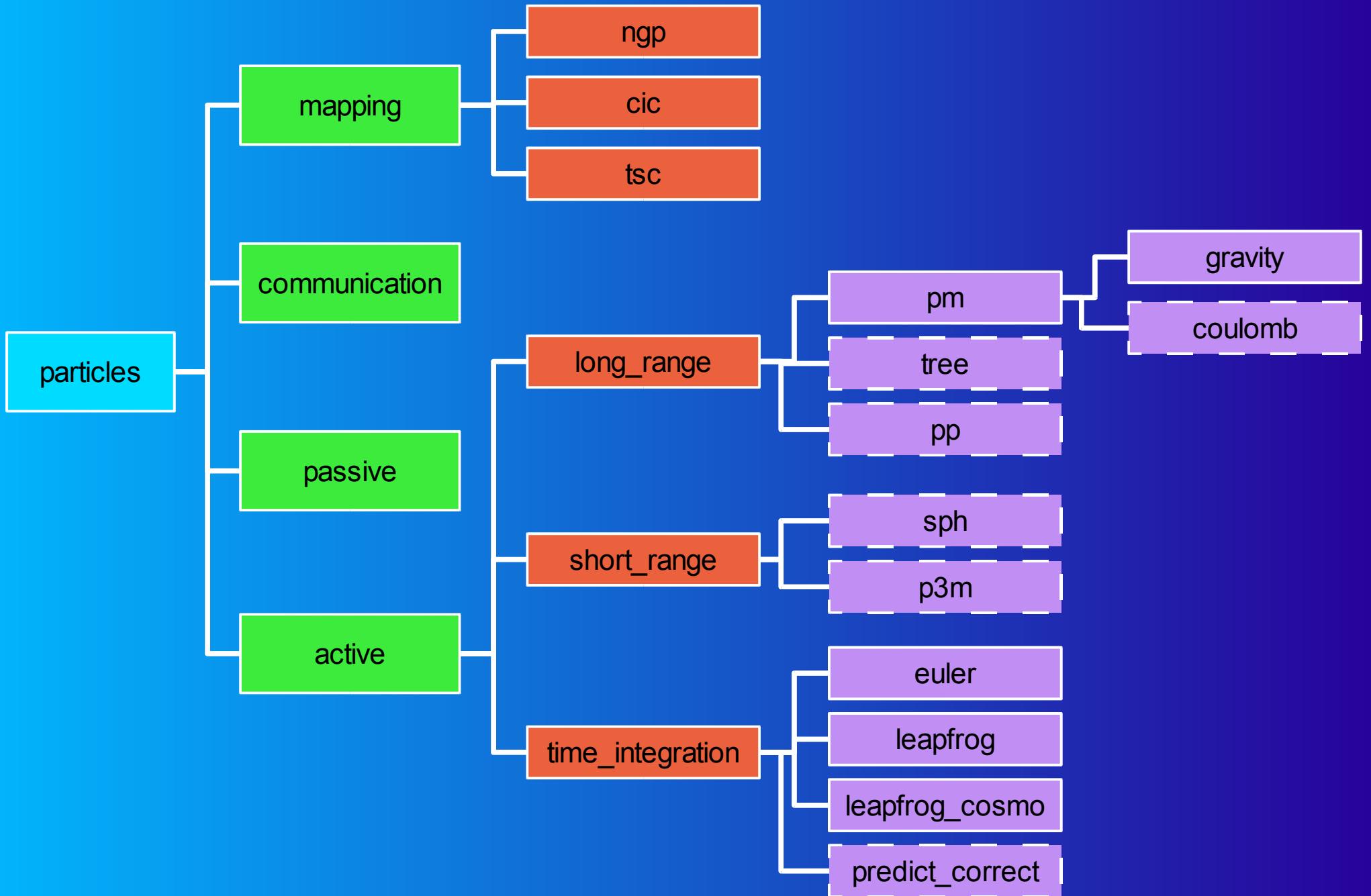


Example: dust cloud collapse



Colgate & White (1966), Mönchmeyer & Müller (1989)

FLASH particle module



Methods supplied by Particles F90 module

● Usage:

```
use ParticleModule, ONLY: <list of methods to import>
```

● Methods:

- `InitParticles()`

Initialize the particle module (called by `init_flash()`)

- `InitParticlePositions(block)`

Initialize particle positions for a given block.

- `AdvanceParticles()`

Time advancement of particle positions.

- `MapMeshToParticles(p_attrib, mesh_var, zero_mode)`

- `MapParticlesToMesh(mesh_var, p_attrib, zero_mode)`

Apply particle-mesh transfer operators.

- `ReDistributeParticles(mesh_mod_flag)`

Distribute particles among processors so that they are colocated with the leaf-node blocks that contain them.

- `ParticleTimestep(dt_part, dt_minloc, block)`

Compute particle timestep restriction.

Data structures supplied by Particles F90 module

- Usage:

```
use ParticleData, ONLY: <list of particle data to import>
```

- Particles not fully integrated with dBase (yet)

- Direct access to particle data structures
- Particle attribute keys managed by dBase

- Basic particle data structure:

```
type particle_type  
    integer :: intAttributes (MaxIntProperties)  
    real      :: realAttributes (MaxRealProperties)  
end type particle_type  
type(particle_type) pointer :: particles(:)
```

- Use RunningParticles (logical) to test if particles are in code and turned on

- pden mesh variable defined for active particles to hold mesh density

Particles FLASH module configuration

```
# Configuration file for the particle module

DEFAULT communication

REQUIRES driver
REQUIRES particles/communication
REQUIRES particles/mapping

EXCLUSIVE active passive

# Parameters:

D  MaxParticlesPerProc      Number of particles to track per processor
D  ipart                      Whether to advance particles or not
D  part_dt_factor             Factor multiplying  $dx/|v|$  in setting particle timestep limit

PARAMETER MaxParticlesPerProc  INTEGER 1
PARAMETER ipart                INTEGER 0
PARAMETER part_dt_factor       REAL    0.5

# Particle properties/attributes

PROPERTY particle_x      REAL
PROPERTY particle_y      REAL
PROPERTY particle_z      REAL
PROPERTY particle_x_vel  REAL
PROPERTY particle_y_vel  REAL
PROPERTY particle_z_vel  REAL
PROPERTY particle_tag    INTEGER
PROPERTY particle_block  INTEGER
```

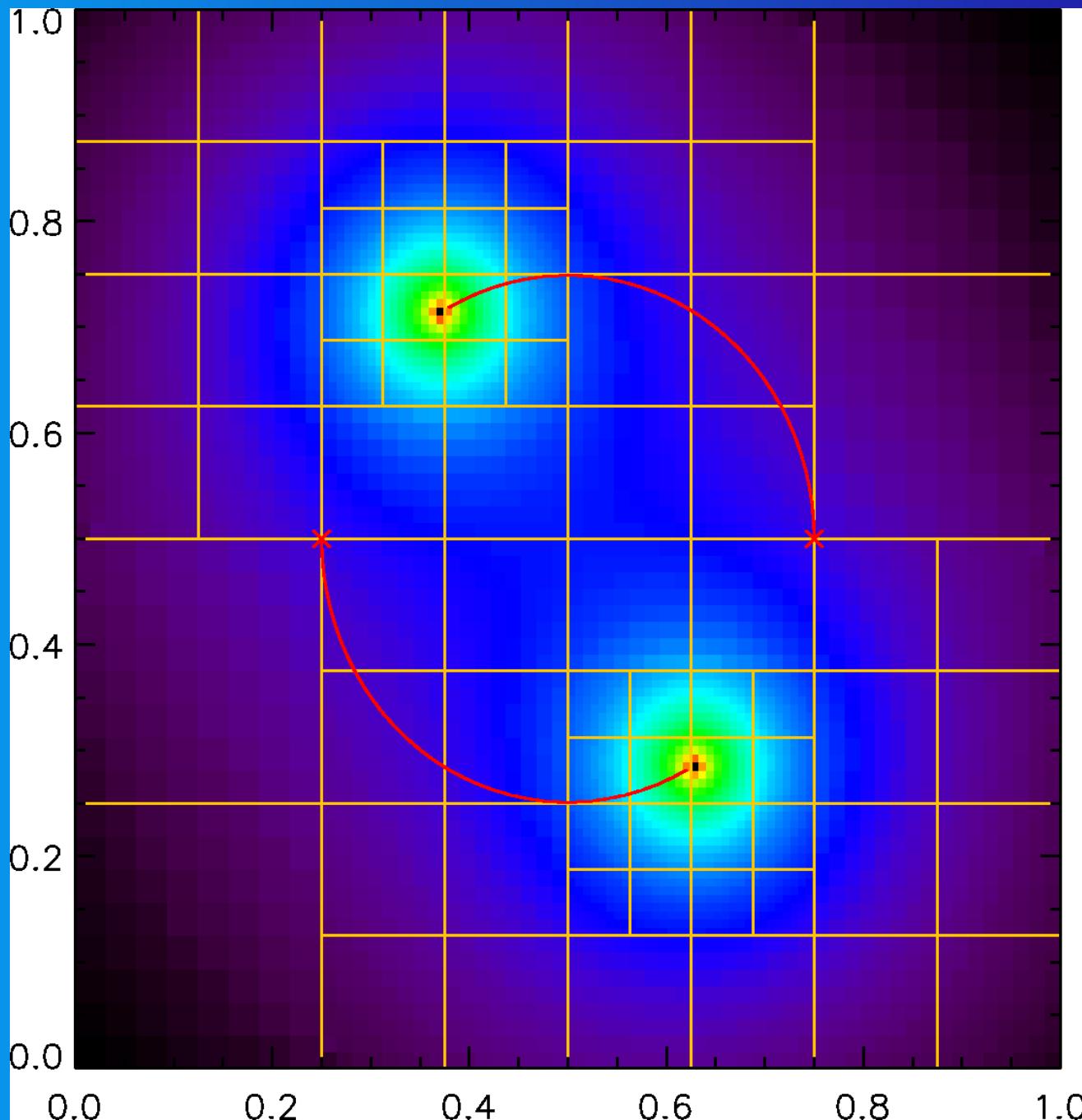
Data structures supplied by Particles F90 module

- Example:

```
use ParticleData, ONLY: particles, ipx, &
                      is_empty
integer :: ipflavor

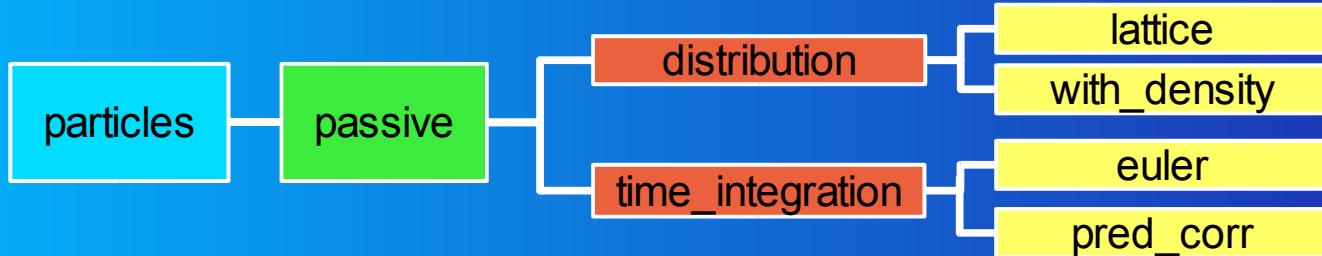
ipflavor = dBaseKey("particle_flavor")
if (.not. is_empty(10)) then
    print *, particles(10)%realAttributes(ipx), &
              particles(10)%realAttributes(ipflavor)
endif
```

Example: orbit problem



Example: tracer particles

- FLASH 2.4 modifies passive particles thusly:

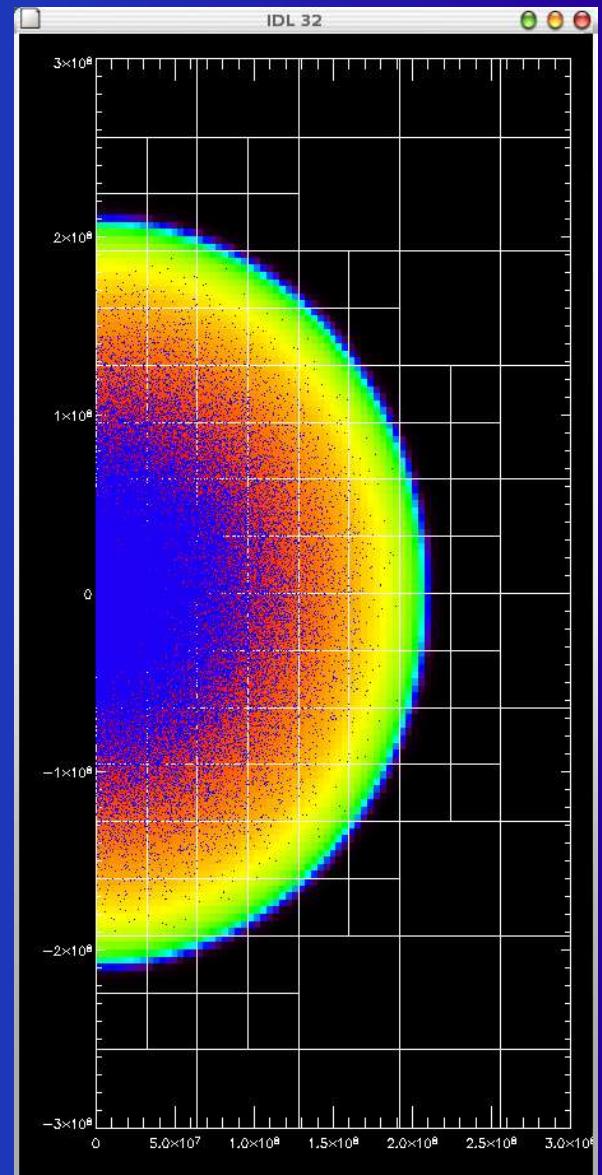


- Use distribution/with_density to distribute particles randomly according to the mesh density

- init_from_scratch() includes

```
use ParticleModule  
... set up mesh using repeated init_block() calls ...  
do block_no = 1, lnblocks  
  if (dbaseNodeType(block_no) == 1) &  
    call InitParticlePositions (block_no)  
  enddo  
  call ReDistributeParticles (.true.)
```

- Use similar approach for particle clouds with specified density profiles (e.g., isothermal sphere)



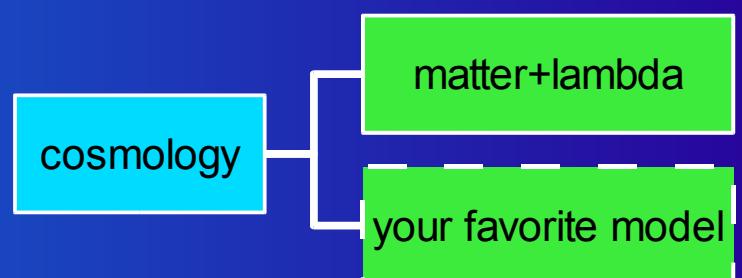
Cosmology FLASH module

```
# Configuration file for the cosmology module.
```

```
DEFAULT matter+lambda
```

```
# Parameters:
```

D OmegaMatter	Ratio of total mass density to closure density at the present epoch
D &	
D OmegaBaryon	Ratio of baryonic mass density to closure density at the present epoch (must be <= OmegaMatter!)
D &	
D CosmologicalConstant	Ratio of the mass density equivalent in the cosmological constant (or dark energy) to the closure density at the present epoch
D &	
D &	
D OmegaRadiation	Ratio of total radiation density to closure density at the present epoch
D &	
D HubbleConstant	Value of the Hubble constant (\dot{a}/a) in sec ⁻¹ at the present epoch
D &	
D MaxScaleChange	Maximum permitted fractional change in the scale factor during each timestep
D &	
PARAMETER OmegaMatter	REAL 0.3
PARAMETER OmegaBaryon	REAL 0.05
PARAMETER CosmologicalConstant	REAL 0.7
PARAMETER OmegaRadiation	REAL 5.E-5
PARAMETER HubbleConstant	REAL 2.1065E-18
PARAMETER MaxScaleChange	REAL 1.E99



Methods supplied by Cosmology F90 module

● Usage:

use Cosmology, ONLY: <*list of methods to import*>

● Methods (sample):

- InitCosmologicalModel () (2.4)

Initialize the cosmology module (called by init_flash())

- SolveFriedmannEquation (t, dt)

Advance scale factor from time t to t+dt.

- RedshiftHydroQuantities ()

Apply operator-split redshift terms in the comoving Euler equations.

- MassToLength (mass, lambda)

Compute comoving diameter of a sphere containing the given mass.

- RedshiftToTime (z, t)

Compute age of the Universe for a given redshift.

- ExpansionTimestep ()

Compute cosmological timestep restriction ($|\Delta a/a| < \text{MaxScaleChange}$).

● Access old and updated scale factor and redshift through dBase

redshift = dBasePropertyReal ("Redshift")

scale = dBasePropertyReal ("ScaleFactor")

Initializing particle problems

- General strategy

- In `init_from_scratch()`, set up mesh (at least approximately)
- Initialize particles once blocks are distributed; use `FindLeafBlockForPosn()` to determine which processor “owns” the particle
- Call `ReDistributeParticles(.true.)` to force particle redistribution

- Some options

- Read particle positions and velocities from a file
- Processor 0 randomly chooses positions from user-supplied or mesh-variable distribution and sends them to processors
- All processors randomly choose positions (rescale # of particles)
- All processors loop through blocks, select known particle positions

Analyzing particle output

- XFLASH (IDL)
 - Can plot particle positions directly
- From command line in IDL using FIDLR routines (FIDLR2)

```
IDL> read_amr_hdf5, "pan1d_hdf5_chk_0000", $  
IDL> particles=p, int_prop_names=inames, $  
IDL> real_prop_names=rnames  
IDL> print, p[10].realAttributes[where(rnames $  
IDL> eq "particle_x")]  
3.5360313e+23
```
- From command line in IDL using FIDLR routines (FIDLR3)

```
IDL> read_amr, "pan1d_hdf5_chk_0000", particles=p  
IDL> print, p[10].particle_x  
3.5360313e+23
```

Analyzing particle output

- From F90/C program using FLASH HDF5 reading routines

```
integer :: file_id, np
integer, parameter :: MAXINT = 12, &
                     MAXREAL = 12 ! must agree with file
type particle_type
    integer :: int(MAXINT)
    real     :: real(MAXREAL)
end type particle_type
type (particle_type) :: p
character(len=24) :: inames(MAXINT), rnames(MAXREAL)

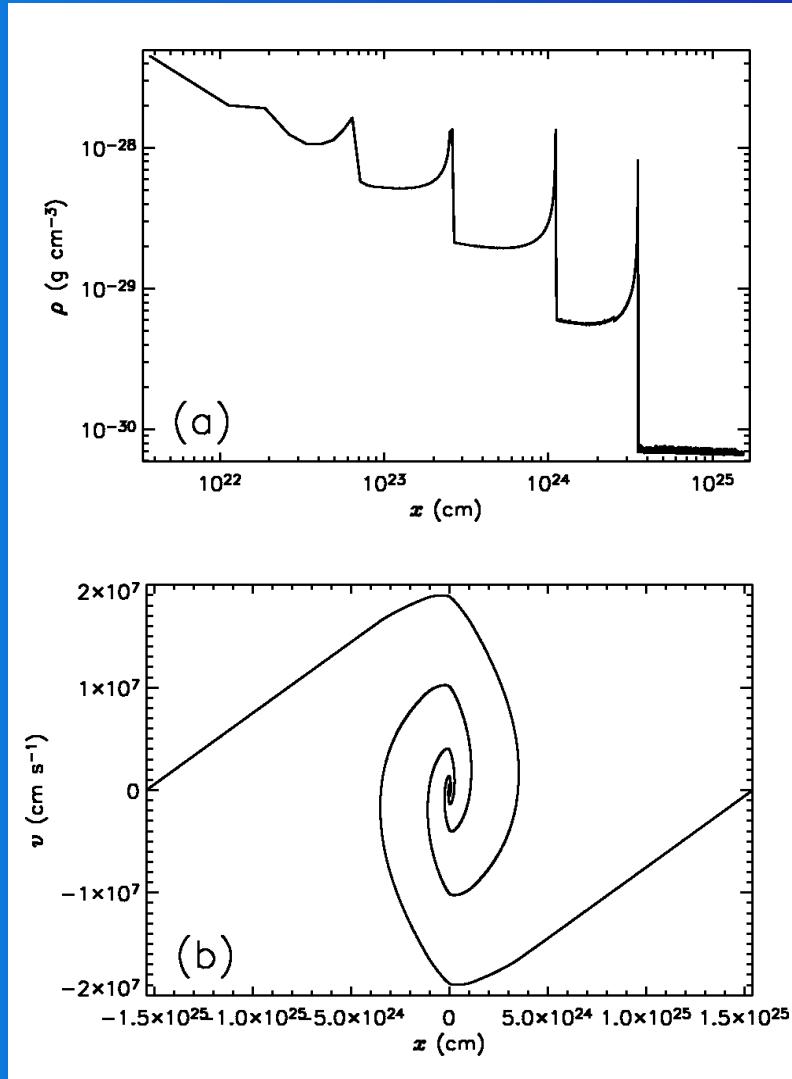
call h5_open_file_for_read (file_id, "myfile")
call get_numparticles (file_id, np)
call h5_read_particles (file_id, np, 1, 9, inames, rnames, p)
print *, p
call h5_close_file (file_id)
```

Example: Zel'dovich pancake

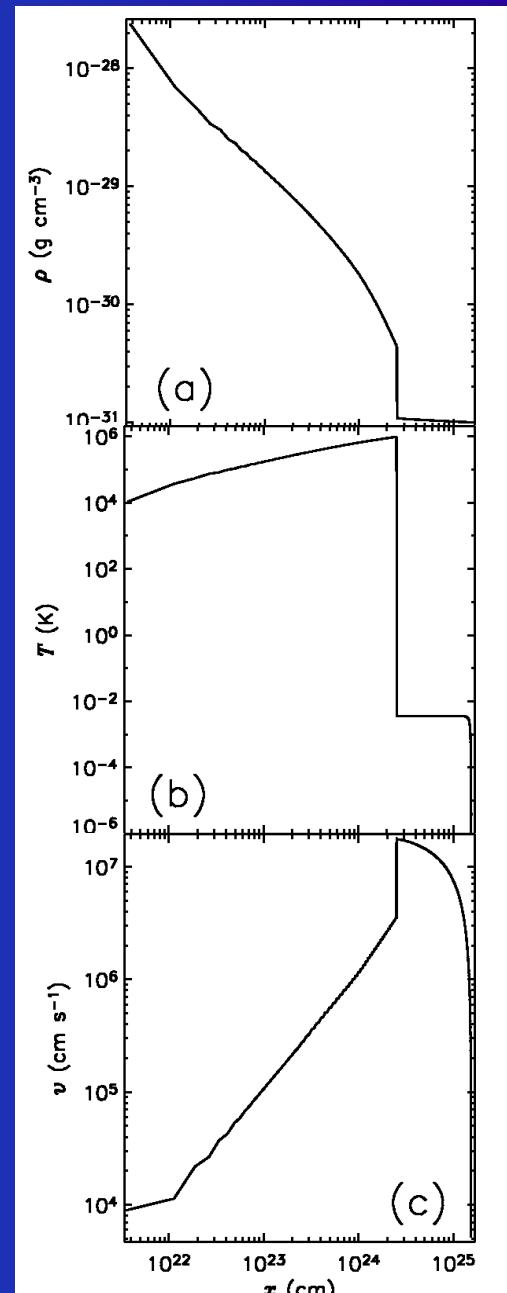
pancake problem

- Single-mode planar collapse
- Collisionless dark matter forms caustics
- Gas forms strong shocks (Mach > 1000)
- Versions:
 - 2.3: create grid of particles, initialize block-by-block
 - 2.4: initialize all at once, discard particles not owned by us

$z = 0$, caustic at $z = 5$
10 levels of refinement



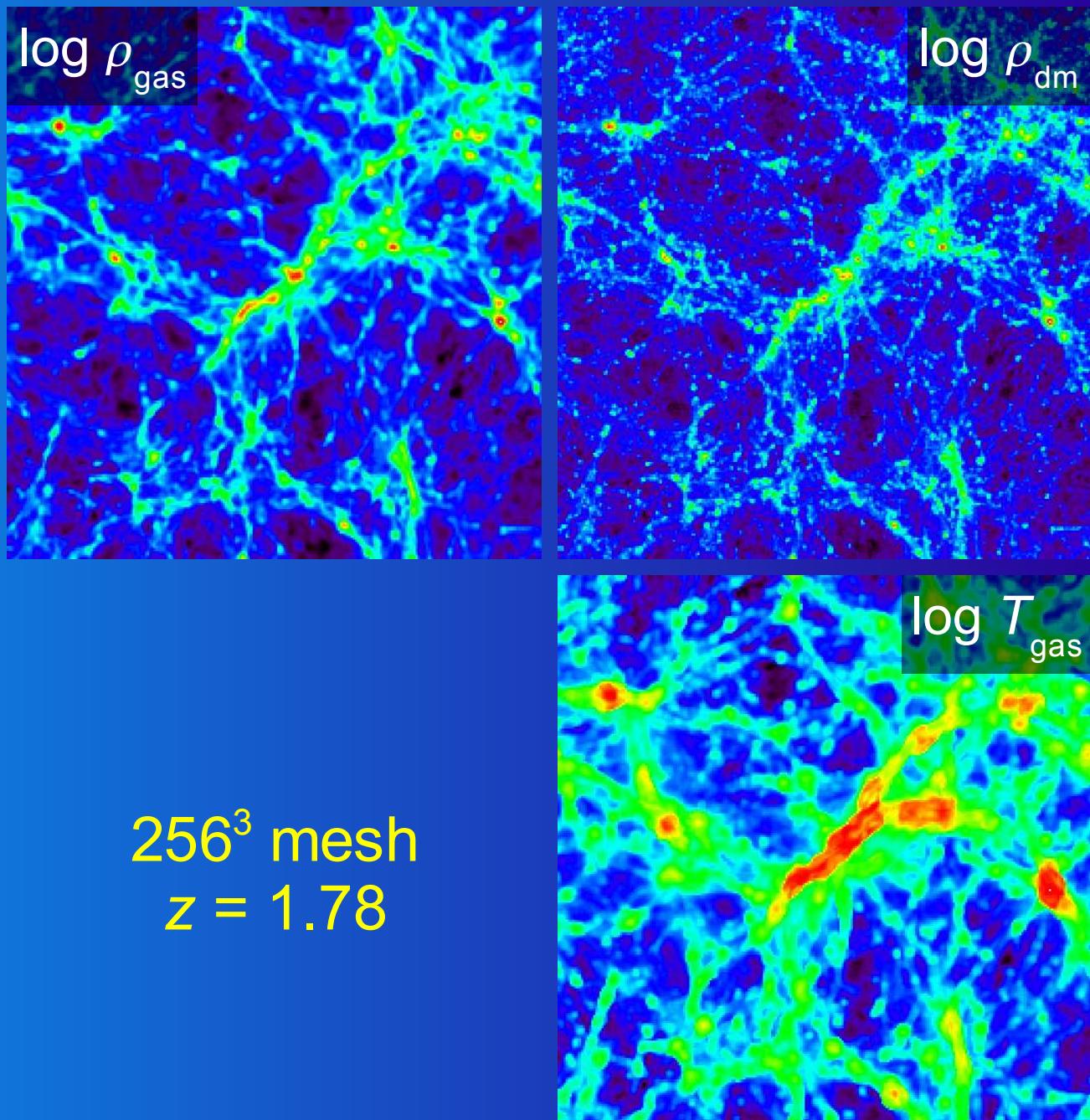
Dark matter



Gas

Example: Santa Barbara cluster

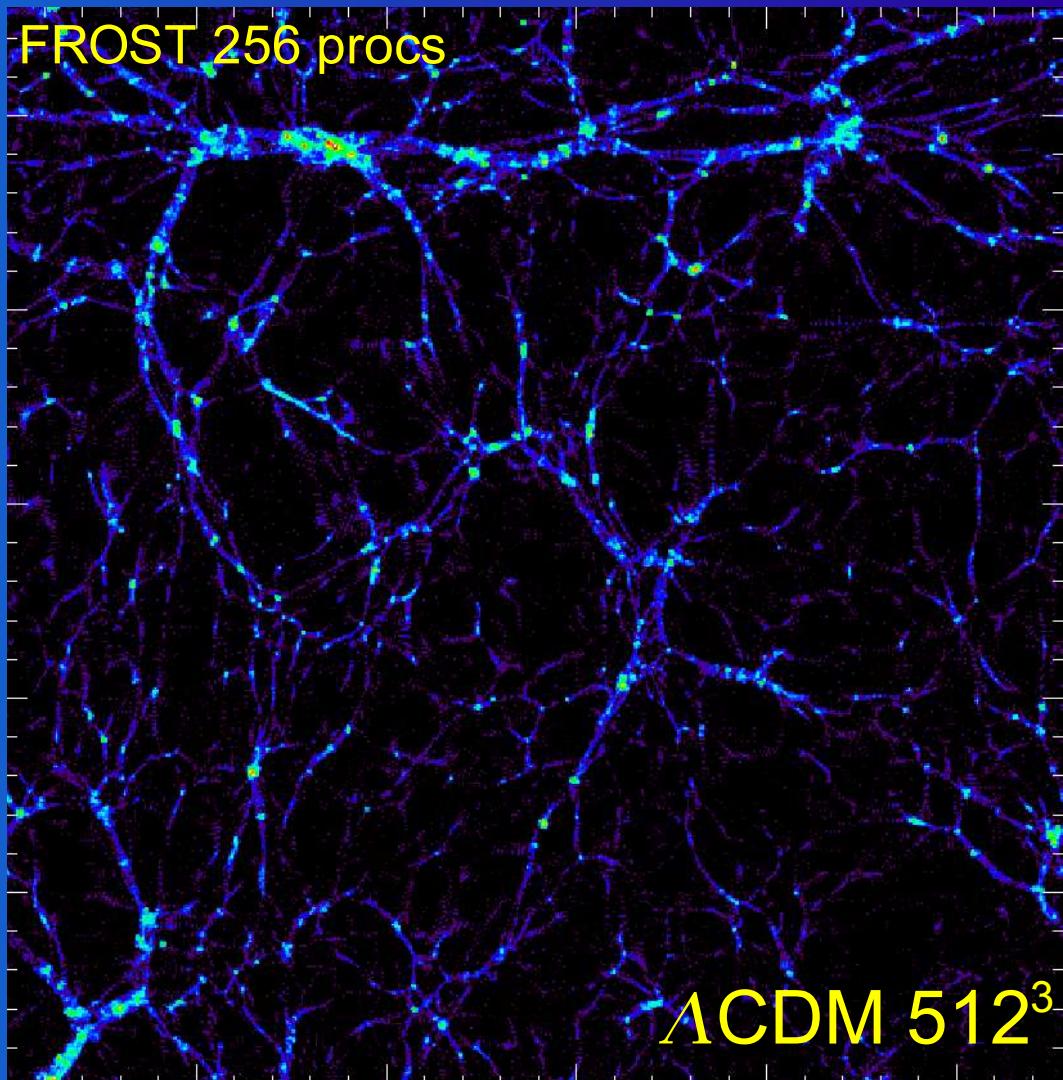
- Standard cold dark matter test problem (Frenk et al. 1999)
- Constrained realization of a $10^{15} M_{\odot}$ cluster
- 256^3 particles
- $(64 \text{ Mpc})^3$ volume
- Initialization from file
- Processor #0 reads, sends particles to owning processors



Example: Λ CDM model

Code comparison with LANL PM code (MC^2), GADGET, TreePM, Warren tree code (HOT) (Heitmann et al. 2004)

- Initialization from file
- Two box sizes
 - $L = 64h^{-1}$ Mpc
 - $L = 256h^{-1}$ Mpc
- 1024^3 particles downsampled to 512^3 and 256^3



Summary

- The big picture
 - Applications for gravity and particles
 - Equations to be solved
- Methods
 - Algorithms currently in FLASH
 - Performance
- Usage
 - Organization of the code modules
 - Initializing a particle application
 - Analyzing particle output
 - Examples