3T Capabilities in FLASH

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RAL Tutorial May 2012 • In the absence of source terms, the hydrodynamic equations are:

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I define a 1T simulation as one in which separate ion,

electron, or radiation energy equations are not evolved

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \boldsymbol{v}) &= 0, \\ \frac{\partial}{\partial t} (\rho \boldsymbol{v}) + \nabla \cdot (\rho \boldsymbol{v} \boldsymbol{v}) + \nabla P_{\text{tot}} &= 0, \\ \frac{\partial}{\partial t} (\rho E_{\text{tot}}) + \nabla \cdot [(\rho E_{\text{tot}} + P_{\text{tot}}) \boldsymbol{v}] &= 0, \end{aligned} \qquad \begin{aligned} E_{\text{tot}} &= e_{\text{ion}} + e_{\text{ele}} + e_{\text{rad}} \\ e_{\text{tot}} &= E_{\text{tot}} - \frac{v^2}{2} \end{aligned}$$

• These equations are closed using an equation of state (EOS):

$$P_{\text{tot}} = \text{EOS}(\rho, e_{\text{tot}})$$

• Since separate ion/electron/radiation energies are not tracked, the 1T EOS units typically make an *assumption* about the temperatures:

$$T_{ele} = T_{ion} = T_{rad}$$
 or $T_{ele} = T_{ion}$, $T_{rad} = 0$

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When is a 1T simulation not adequate?



- Source terms may preferentially heat electrons, ions, or the radiation field creating different temperatures
 - Shocks preferentially heat ions
 - Laser preferentially heats electrons
- The difference in temperatures may be important for:
 - Evaluating transport coefficients (conductivity, resistivity, ...)
 - Evaluating the EOS
- Thus, it is necessary to evolve the electron, ion, and radiation state so that we can keep track of the different electron/ion/radiation internal energies, temperatures, etc...
- All laser driven HEDP experiments should use the 3T version of FLASH, to activate just specify the following setup option:
 - Split Hydrodynamics Solver: +3t
 - Unsplit Hydrodynamics Solver: +uhd3t
 - Staggered Mesh Unsplit MHD Solver: +usm3t



• The same 1T hydrodynamic equations:

$$\begin{split} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \boldsymbol{v}) &= 0, \\ \frac{\partial}{\partial t} (\rho \boldsymbol{v}) + \nabla \cdot (\rho \boldsymbol{v} \boldsymbol{v}) + \nabla P_{\text{tot}} &= 0, \\ \frac{\partial}{\partial t} (\rho E_{\text{tot}}) + \nabla \cdot [(\rho E_{\text{tot}} + P_{\text{tot}}) \, \boldsymbol{v}] &= 0, \\ e_{\text{tot}} &= E_{\text{tot}} - \frac{v^2}{2} \end{split}$$

• Separate equations for ion, electron, radiation internal energy:

$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,{\rm ele}}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,{\rm ele}}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis} + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} &= \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$
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$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis} + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} = \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$

• These equations are closed using an EOS, but now, information about the electron, ion, radiation state are known, so the EOS can use this information:

$$P_{\text{tot}}, T_{\text{ion}}, T_{\text{ele}}, T_{\text{rad}} = \text{EOS}(\rho, e_{\text{ion}}, e_{\text{ele}}, e_{\text{rad}})$$
$$P_{\text{tot}}, e_{\text{ion}}, e_{\text{ele}}, e_{\text{rad}} = \text{EOS}(\rho, T_{\text{ion}}, T_{\text{ele}}, T_{\text{rad}})$$

 Typically, for simulations of HEDP experiments, EOS tables are used for each material



$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis} + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} &= \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$

- $q_{qele} = -K_{ele} \nabla T_{ele}$ represents the electron heat flux, where K_{ele} is the electron conductivity
- A flux-limiter is used to limit the electron heat flux in regions where the temperature gradient is steep
- The Spitzer conductivity is primarily used in FLASH simulations, although the next release will include the Lee-More conductivity

Electron thermal conduction (q_{ele}) is advanced in the Diffuse unit. It is solved implicitly using the HYPRE parallel linear algebra library



• To access the implicit HYPRE based diffusion solver include the following line in your Config file:

REQUESTS physics/Diffuse/DiffuseMain/Unsplit

• FLASH solves the following equation to model electron thermal conduction:

$$\rho c_{v,\text{ele}} \frac{dT_{\text{ele}}}{dt} = \nabla \cdot K_{\text{ele}} \nabla T_{\text{ele}}^*$$

• Typically the Sptizer conductivity is used. To access it, include:

REQUESTS physics/materialProperties/Conductivity/ConductivityMain/SpitzerHighZ

$$K_{\rm ele} = \left(\frac{8}{\pi}\right)^{3/2} \frac{k_B^{7/2}}{e^4 \sqrt{m_{\rm ele}}} \left(\frac{1}{1+3.3/\bar{z}}\right) \frac{T_{\rm ele}^{5/2}}{\bar{z} \ln \Lambda_{ei}}$$



$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis} + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} = \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$

- The ion/electron equilibration term causes the electron/ion temperatures to relax over time due to collisions
- τ_{e} is the ion/electron collision frequency. Most simulations use the Spizter form for this term

Ion/electron equilibration is advanced by the HeatExchange unit



• To access ion/electron equilibration using the Spitzer collision time, include:

REQUESTS physics/sourceTerms/Heatexchange/HeatexchangeMain/Spitzer

• FLASH solves the following equations to model ion/electron equilibration:

$$\frac{\partial e_{\rm ion}}{\partial t} = \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}),$$
$$\frac{\partial e_{\rm ele}}{\partial t} = \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}).$$

• The ion/electron equilibration time is:

$$\tau_{ei} = \frac{3k_B^{3/2}}{8\sqrt{2\pi}e^4} \frac{(m_{\rm ion}T_{\rm ele} + m_{\rm ele}T_{\rm ion})^{3/2}}{(m_{\rm ele}m_{\rm ion})^{1/2} \bar{z}^2 n_{\rm ion}\ln\Lambda_{ei}}$$



$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis} + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} = \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$

- Q_{las} represents the energy source due to laser energy deposition. This is computed using a ray-tracing algorithm (this will be discussed in detail tomorrow)
- The user can specify the location of multiple beams which illuminate the target
- The laser energy is absorbed using the common inverse Bremsstrahlung absorption coefficient



$$\begin{split} \frac{\partial}{\partial t}(\rho e_{\rm ion}) + \nabla \cdot (\rho e_{\rm ion} \boldsymbol{v}) + P_{\rm ion} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ele} - T_{\rm ion}), \\ \frac{\partial}{\partial t}(\rho e_{\rm ele}) + \nabla \cdot (\rho e_{\rm ele} \boldsymbol{v}) + P_{\rm ele} \nabla \cdot \boldsymbol{v} &= \rho \frac{c_{v,\rm ele}}{\tau_{ei}} (T_{\rm ion} - T_{\rm ele}) - \nabla \cdot \boldsymbol{q}_{\rm ele} + Q_{\rm abs} - Q_{\rm emis}) + Q_{\rm las}, \\ \frac{\partial}{\partial t} (\rho e_{\rm rad}) + \nabla \cdot (\rho e_{\rm rad} \boldsymbol{v}) + P_{\rm rad} \nabla \cdot \boldsymbol{v} = \nabla \cdot \boldsymbol{q}_{\rm rad} - Q_{\rm abs} + Q_{\rm emis}, \end{split}$$

- The third equation above represents the change in the radiation energy density where q_{rad} is the radiation energy flux, Q_{emis} is the radiation energy emission, and Q_{abs} is the absorption of energy
- The flux is computed using flux-limited multigroup diffusion with tabulated opacities. This will be discussed in detail tomorrow

FLASH solves the 3T equations by operator splitting the source terms from the "hydrodynamic" terms. Each source term is advanced by a different code unit...



- Electron thermal conduction (q_{ele}) is advanced in the Diffuse unit. It is solved implicitly using the HYPRE parallel linear algebra library
- Laser energy deposition is computed in the EnergyDeposition unit
- Radiation emission, absorption, and energy flux is advanced in the RadTrans unit. This is also evaluated implicitly using HYPRE
- Ion/electron equilibration is advanced by the HeatExchange unit
- The Eos unit computes the 3T equation of state
- The remaining terms in the energy equations represent hydrodynamic work and advection of internal energy by the fluid; these effects are included by the Hydro unit

FLASH solves the 3T equations by operator splitting the source terms from the "hydrodynamic" terms, we are left with:



$$\begin{split} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \boldsymbol{v}) &= 0, \\ \frac{\partial}{\partial t} (\rho \boldsymbol{v}) + \nabla \cdot (\rho \boldsymbol{v} \boldsymbol{v}) + \nabla P_{\text{tot}} &= 0, \\ \frac{\partial}{\partial t} (\rho E_{\text{tot}}) + \nabla \cdot [(\rho E + P_{\text{tot}}) \, \boldsymbol{v}] &= 0, \\ \frac{\partial}{\partial t} (\rho e_{\text{ion}}) + \nabla \cdot (\rho e_{\text{ion}} \boldsymbol{v}) + P_{\text{ion}} \nabla \cdot \boldsymbol{v} &= 0, \\ \frac{\partial}{\partial t} (\rho e_{\text{ele}}) + \nabla \cdot (\rho e_{\text{ele}} \boldsymbol{v}) + P_{\text{ele}} \nabla \cdot \boldsymbol{v} &= 0, \\ \frac{\partial}{\partial t} (\rho e_{\text{rad}}) + \nabla \cdot (\rho e_{\text{rad}} \boldsymbol{v}) + P_{\text{rad}} \nabla \cdot \boldsymbol{v} &= 0, \end{split}$$

- Solving these equations can be difficult for an Eulerian code, the work terms are NOT written as the divergence of a flux
- Worse, they include the divergence of velocity, which is discontinuous at a shock

The electron, ion, radiation, internal energy equations cannot be directly evaluated, a work around is needed...



- In general, there are three hydrodynamic effects which alter the ion, electron, radiation, and total internal energy:
 - Advection
 - Work
 - Shock heating
- The normal 1T hydrodynamics solver updates the total internal energy. The change in internal energy in a given cell must then be divided among the ions, electrons, and radiation field
- For example, the total work done on a cell (or by a cell) is divided among the ions, electrons, and radiation field in proportion to their pressures

One work around involves treating the electrons adiabatically (no shock heating for electrons)



• To do this, you replace the electron internal energy equation with an advection equation for the electron entropy

$$\begin{split} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \boldsymbol{v}) &= 0, \\ \frac{\partial}{\partial t} (\rho \boldsymbol{v}) + \nabla \cdot (\rho \boldsymbol{v} \boldsymbol{v}) + \nabla P_{\text{tot}} = 0, \\ \frac{\partial}{\partial t} (\rho E_{\text{tot}}) + \nabla \cdot [(\rho E + P_{\text{tot}}) \, \boldsymbol{v}] &= 0, \\ \frac{\partial}{\partial t} (\rho s_{\text{ele}}) + \nabla \cdot (\rho s_{\text{ele}} \boldsymbol{v}) &= 0, \\ e_{\text{tot}} &= E_{\text{tot}} - \frac{v^2}{2}, \\ P_{\text{tot}}, e_{\text{ion}}, T_{\text{ele}}, T_{\text{ion}}, e_{\text{ele}}, \dots &= \text{EOS}(\rho, e_{\text{tot}}, s_{\text{ele}}) \end{split}$$

- Unfortunately, the electron entropy method currently only works without radiation and for a fixed-ionization gamma-law EOS
- But we (Klaus) are working on extending this so that it works generally!

Structure of a shock in a plasma with electron thermal conduction and no radiation





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The "RAGE-like" hydro method is an alternative to entropy-advection, it is less accurate near shocks



- The entropy-advection approach treats the electrons adiabatically
- An alternative is to divide the shock heating between electrons in exactly the same way as the work in proportion to the pressure ratios
 - This is the "RAGE-like" approach to 3T hydrodynamics in an Eulerian code
- In this case, the electrons would be shock heated, but the solution away from shocks should be accurate
- This is not necessarily important in plasmas the ion/electron temperatures equilibrate over time!
- RAGE-like hydro is the default method, you can try entropy advection by setting the following parameter in the runtime parameters file (flash.par):

```
hy_eosModeAfter = "dens_ie_sele_gather"
```

There is a large difference between entropy-advection and RAGE-like hydro in the absence of ion/electron equilibration





When Ion/electron equilibration is significant, the difference between RAGE-like hydro and entropy advection is small – even near a shock





FLASH has been extended to support 3T simulations



- 3T should be used for simulations of laser driven HEDP experiments
- A lot of physics relevant to simulating HEDP experiments only works with 3T:
 - Laser
 - Electron thermal conduction
 - Multrigroup radiation diffusion
- 3T hydrodynamics in an Eulerian code is a little tricky because the electrons should be treated adiabatically
- This can affect the electron/ion temperatures near shocks
- We will see some hands on examples and a details on the laser and multigroup diffusion radiation tomorrow!
- For more on 3T simulations, please see user's guide (Ch 13, Sec 25.7.5):

http://flash.uchicago.edu/site/flashcode/user_support/flash4b_ug/node20.html